



QUit
quantum information
theory group

PHYSICS AS INFORMATION THEORY

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OUTLINE

- * Informational axiomatization of Quantum Theory
- * How space-time and relativistic covariance emerge from the quantum computation
- * What is the information-theoretical meaning of inertial mass and \hbar , and how the quantum field emerges
- * Observational consequences: mass-dependent refraction index of vacuum

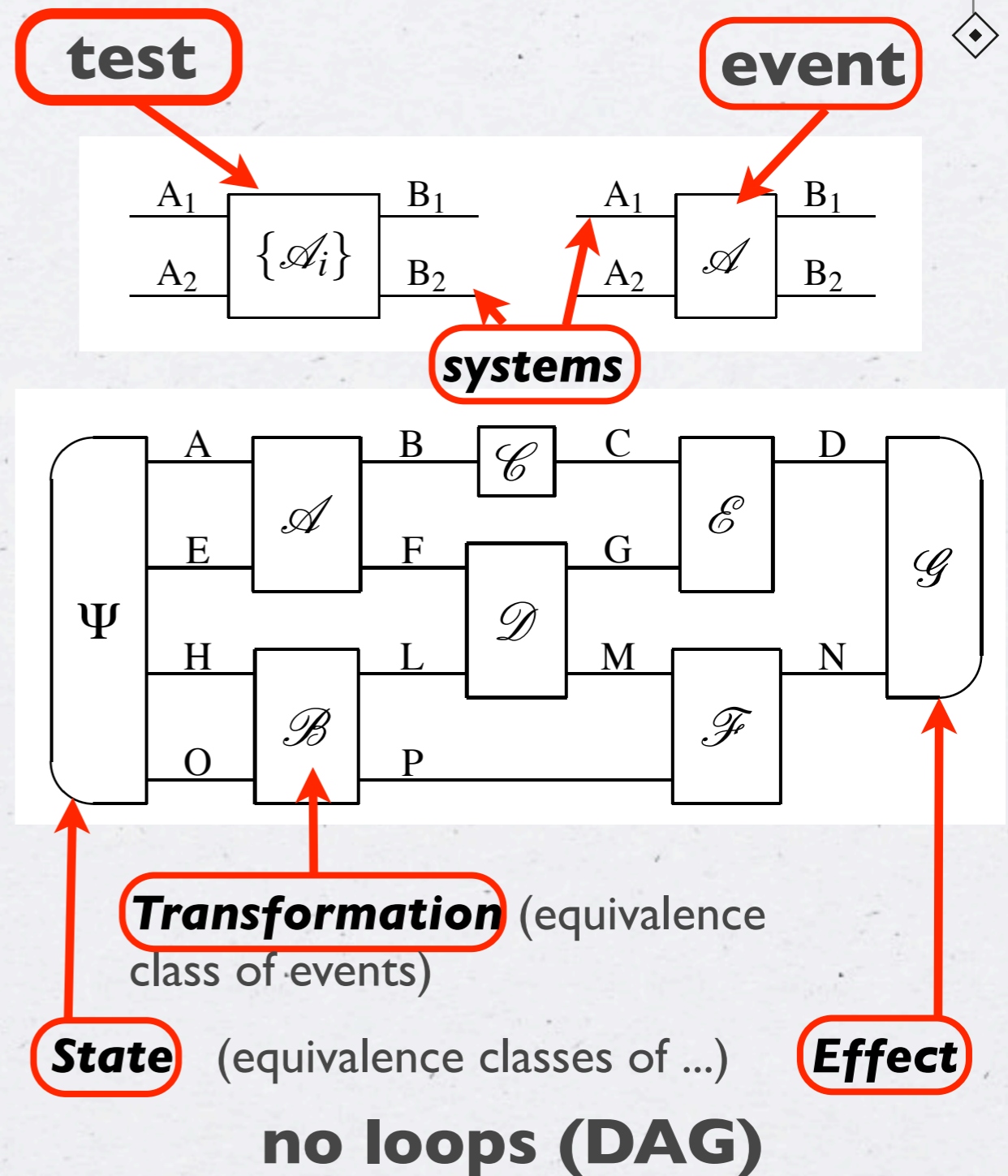
OPERATIONAL FRAMEWORK

Notions: coarse-graining, refinement, refinement set, atomic/indivisible

Probabilistic operational theory: every closed circuit made of events is associated to a probability.

D'Ariano in *Philosophy of Quantum Information and Entanglement*, A. Bokulich and G. Jaeger (Cambridge Un. Press 2010)

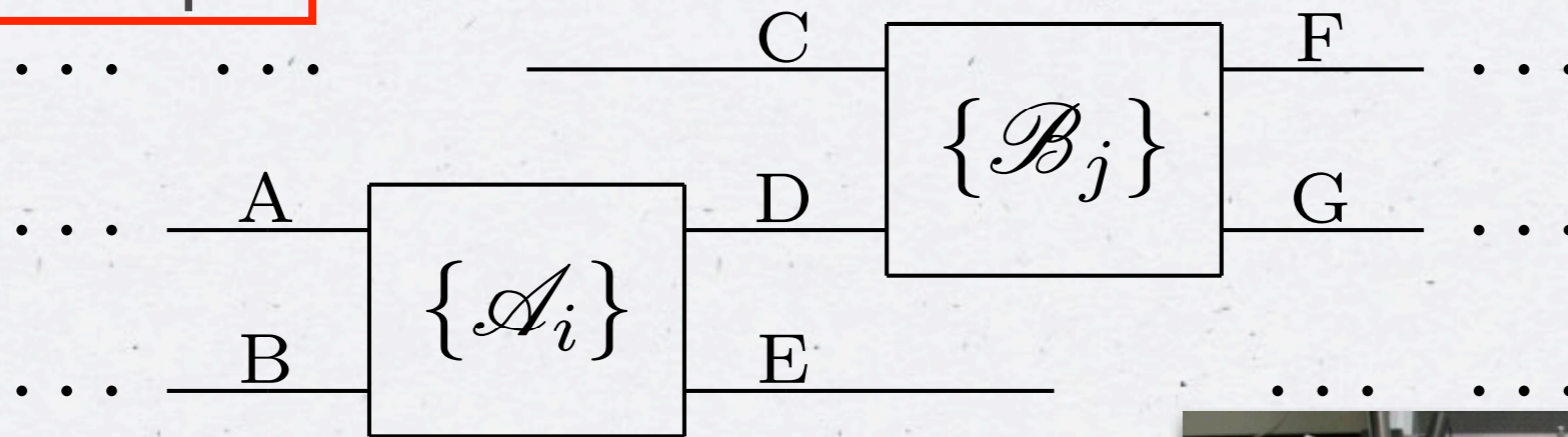
Chiribella, D'Ariano, and Perinotti, PRA **81** 062348 (2010)



CAUSAL PROBABILISTIC THEORIES

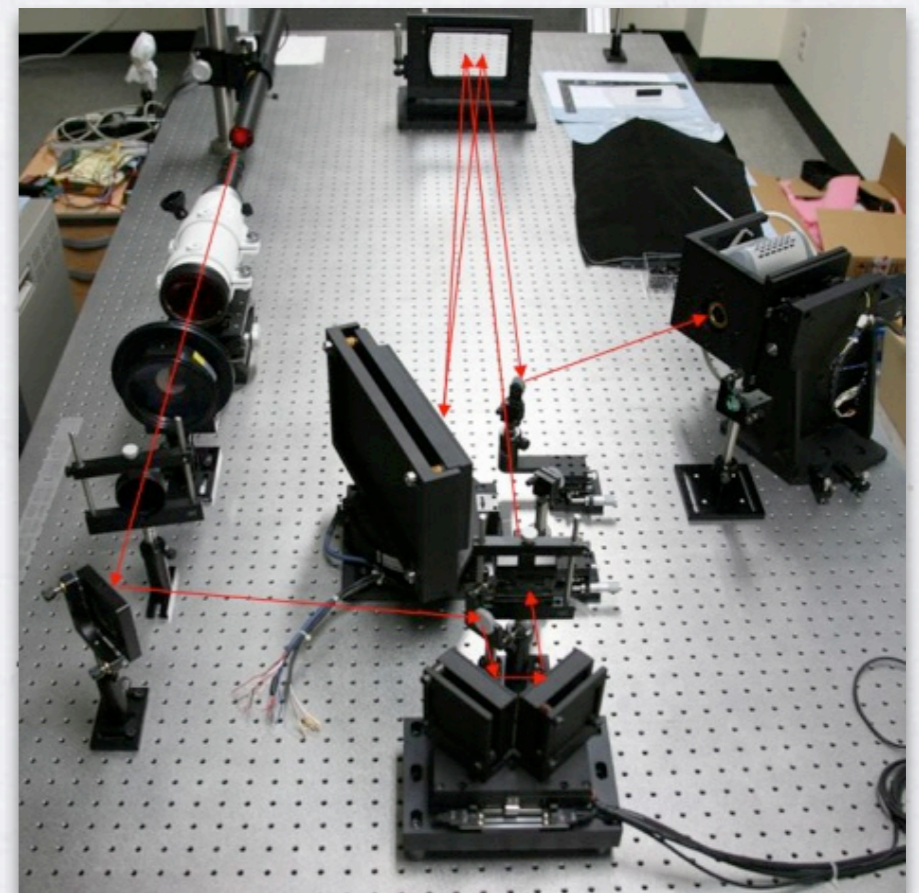
Input \rightarrow Output

DAG



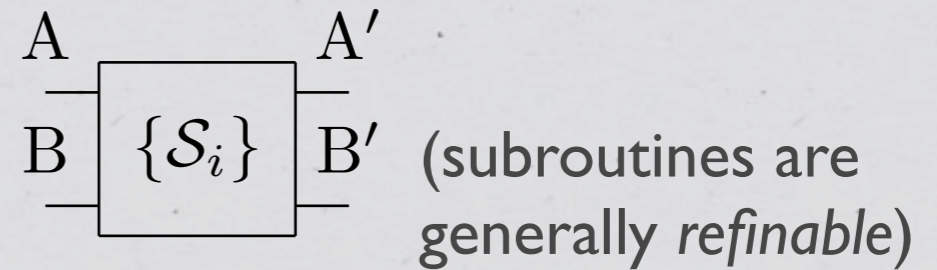
A theory is *causal* if for any two tests that are input-output connected the marginal probability of the input event is independent on the choice of the output test.

G. M. D'Ariano in *Philosophy of Quantum Information and Entanglement*, A. Bokulich and G. Jaeger (CUP, Cambridge UK, 2010).

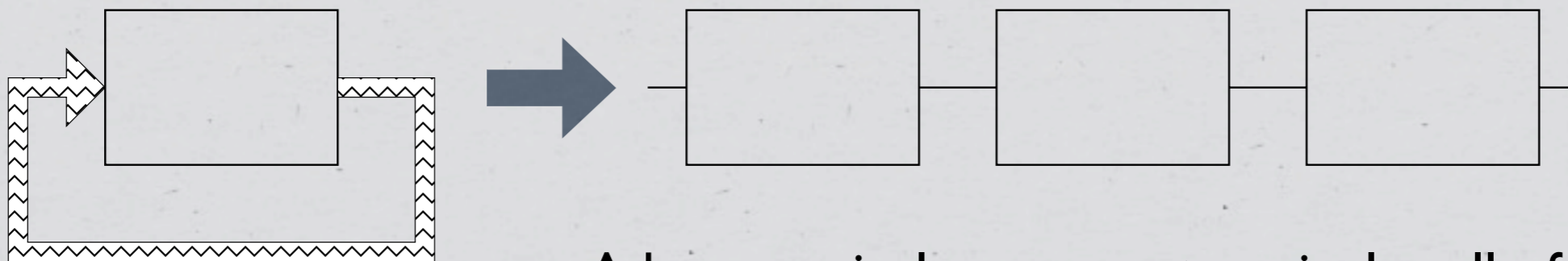
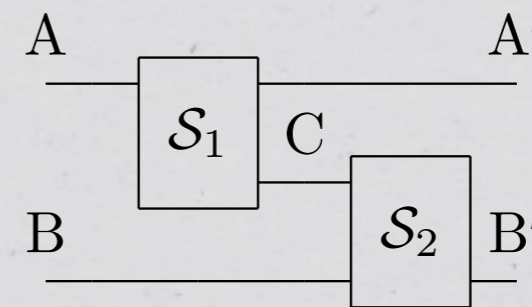


TRANSLATION IN TERMS OF INFORMATION PROCESSING

- *Test/Event* → **subroutine**
- *Transformation* → **information processing**
- *System* → **register**
- *States* → **initialization**
- *Effects* → **readout**
- ... *Pure state* → **indivisible initialization, etc..**



We can **compose processings** connecting input with outputs of the same type



- A box precisely represents a single call of the processing
- The circuit represents the entire run, not a flow diagram

TRANSLATION IN TERMS OF INFORMATION PROCESSING

$$\boxed{A} \xrightarrow{A} \boxed{S} \xrightarrow{B} = \boxed{A'} \xrightarrow{B} \quad \text{Initialization followed by a processing} = \text{new initialization}$$

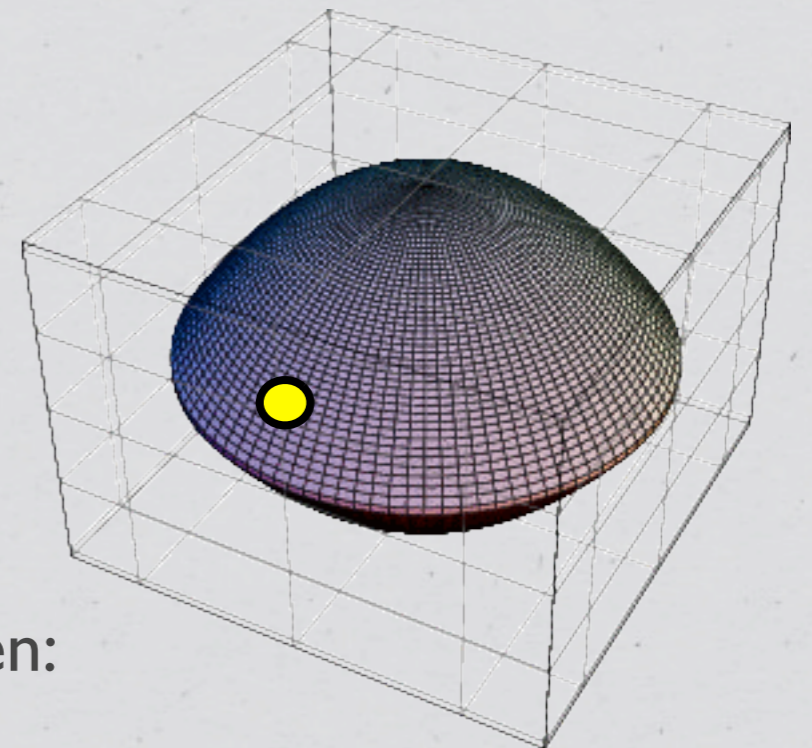
The **domain** of a processing is the set of its possible initializations, its **range** the set of its possible readouts.

An initialization is **specific** when its refinement set is not the whole set of initializations.

Two initializations \mathcal{A}_1 and \mathcal{A}_2 are **discriminable** when:

$$\boxed{\mathcal{A}_1} \xrightarrow{A} \boxed{B} \neq \boxed{\mathcal{A}_2} \xrightarrow{A} \boxed{B}$$

and the discrimination is **perfect** when the two probabilities are 0 and 1



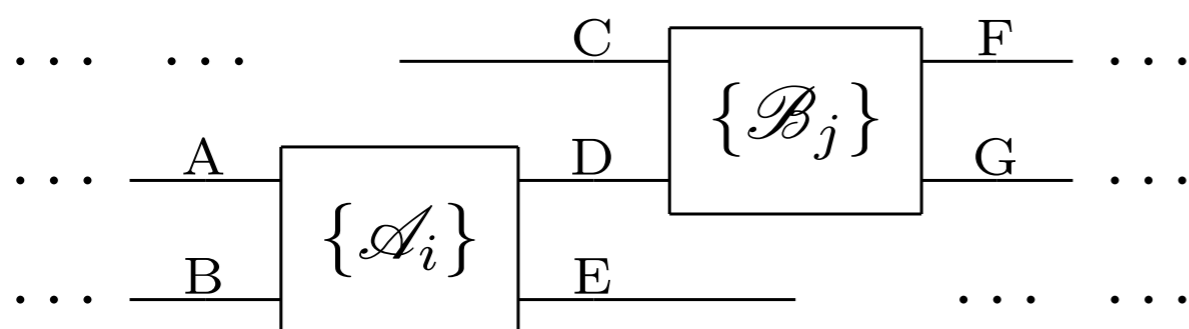
THE PRINCIPLE OF THE QUANTUMNESS

POSTULATES FOR QT

P1. **Causality:** The occurrence of a component processing cannot depend on the choice of the processing at its output (i. e. information flows only from input to output).

P2. **Local Readability:** We can discriminate two initializations of multiple registers by readouts on single registers.

P3. **Reversibility and Indivisibility of Computation:** Every information processing can be replaced with a reversible one by adding a register in an indivisible



P4. **Indivisibility of Processing:** The processing corresponding to a specific initialization is itself indivisible.

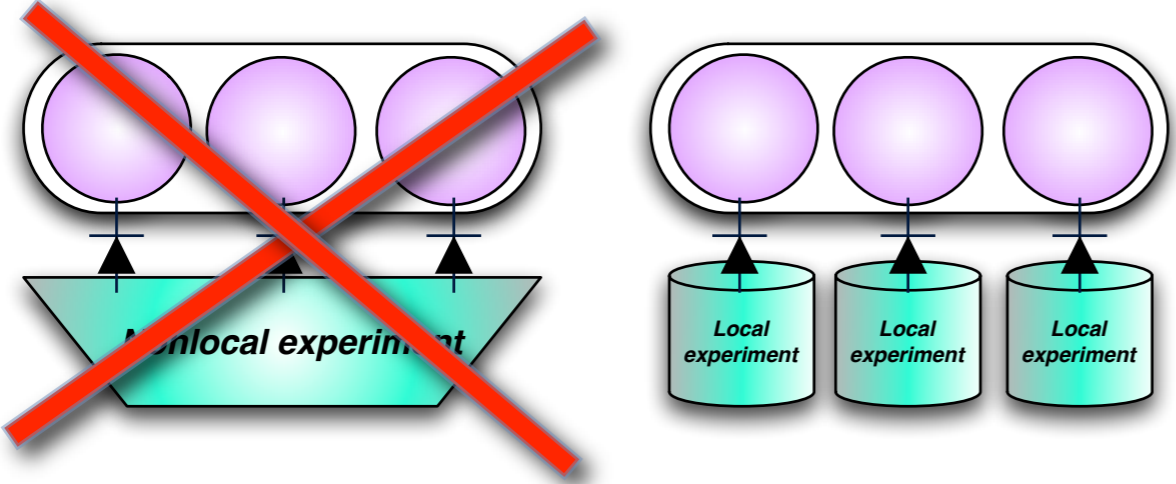
P5. **Discriminability of Specific Initializations:** For any specific initialization there exists another initialization that can be perfectly discriminated from it.

P6. **Lossless Compressibility:** For any initialization there exists an encoding which is perfectly decodeable on its refinement set, and the encoded initialization is not specific.

THE

THE ONLY

SS



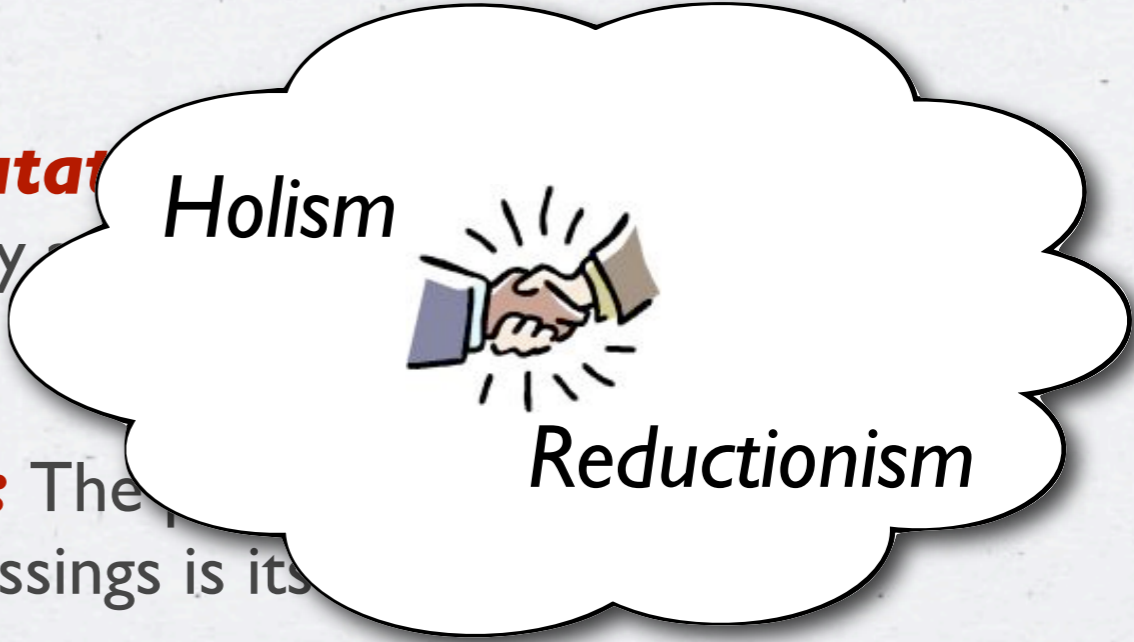
THE ORIGIN OF THE COMPLEX TENSOR PRODUCT

★ ★ ★

P1. **Choice** are processing at its output (i.e. information flows only from input to output).

P2. **Local Readability:** We can discriminate two initializations of multiple registers by readouts on single registers.

P3. **Reversibility and Indivisibility of Computation** processing can be achieved with a reversible one by initialization.



CRUCIAL IN REDUCING EXPERIMENTAL COMPLEXITY, BY GUARANTEEING THAT ONLY LOCAL (JOINTLY EXECUTED) MEASUREMENTS ARE SUFFICIENT TO RETRIEVE A COMPLETE INFORMATION OF A COMPOSITE SYSTEM, INCLUDING ALL CORRELATIONS BETWEEN THE COMPONENTS

P6. **Lossless Compressibility:** For any initialization there exists an encoding which is perfectly decodeable on its refinement set, and the encoded initialization is not specific.

THE PRINCIPLE OF THE QUANTUMNESS

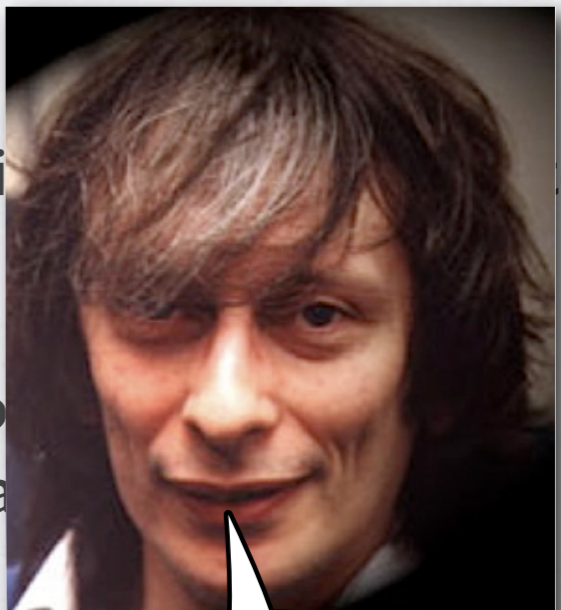
POSTULATES FOR QT

P1. **Causality:** The occurrence of a component processing cannot depend on the choice of the processing at its output (i. e. information flows only from input to output).

P2. **Local Readability:** We can discriminate two initializations of multiple registers by readouts on single registers.

P3. **Reversibility and Indivisibility of Computation:** Every information processing can be achieved with a reversible one by adding a register in an indivisible initialization.

P4. **Indivisibility of Processing Components:** For any initial state of two indivisible components, there exists a processing component that can be achieved with a reversible one by adding a register in an indivisible initialization.



COMPUTATION CAN BE DONE REVERSIBLY

COMPUTATION RUNS IN QUANTUM PARALLELISM

THE MOST "QUANTUM" POSTULATE

- ALL POSTULATES APART FROM P3 ARE SATISFIED BY CLASSICAL THEORY, P3 IS NOT SATISFIED BY PR BOXES
- NO KNOWN THEORY (APART FROM QT) SATISFYING P1, P2, AND P3
- IT IS THE BASIS OF MOST QUANTUM INFORMATION PROTOCOLS: TELEPORTATION, ERROR CORRECTION, NO-CLONING THEOREM, ANCILLA-ASSISTED TOMOGRAPHY,

THE PRINCIPLE OF THE QUANTUMNESS

POSTULATES FOR QT

- ★ ★ ★ ★
- P1. **Causality:** The information processing cannot depend on the choice of the initialization (NO REASON WHY THE SAME INFORMATION FLOWS ONLY FROM INPUT TO OUTPUT).
PROCESSING OBTAINED BY COMPOSING TWO ONES COULD NOT BE ITSELF ACHIEVED IN PRINCIPLE BY A SUBROUTINE
- P2. **Local Readability:** The information processing cannot be achieved by readouts of multiple registers without initializations of multiple registers.
- P3. **Reversibility:** The information processing can be achieved by initialization of multiple registers. WHICH IS DIVISIBLE
- P4. **Indivisibility of Processing Composition:** The processing corresponding to the input-output sequence of two indivisible processings is itself indivisible.
- P5. **Discriminability of Specific Initializations:** For any specific initialization there exists another initialization that can be perfectly discriminated from it.
- P6. **Lossless Compressibility:** For any initialization there exists an encoding which is perfectly decodeable on its refinement set, and the encoded initialization is not specific.
- ★ ★ ★ ★

THE PRINCIPLE OF THE QUANTUMNESS

POSTULATES FOR QT

★ ★ ★ ★

P1. **Causality:** The occurrence of a component processing cannot depend on the choice of the processing at its output (i. e. information flows only from input to output).

P2. **Local Readability:** We can discriminate two initializations of multiple registers by readouts on

P3. **Reversibility:** Every information processing can be made reversible one by adding a register in an indivisible initialization.

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P6. **Lossless Compressibility:** For any initialization there exists an encoding which is perfectly decodeable on its refinement set, and the encoded initialization is not specific.

IT IS EASY TO
CONSTRUCT A
THEORY THAT
VIOLATES IT

THE PRINCIPLE OF THE QUANTUMNESS

POSTULATES FOR QT



P1. **Computational Irreducibility:** Processing cannot depend on the state of the system (i.e., it always only from input to output). This is NOT OBVIOUS FOR INFORMATION PROCESSING WITH DIFFERENT KINDS OF REGISTERS. It is CRUCIAL FOR SHANNON'S & SCHUMAKER'S COMPRESSION.

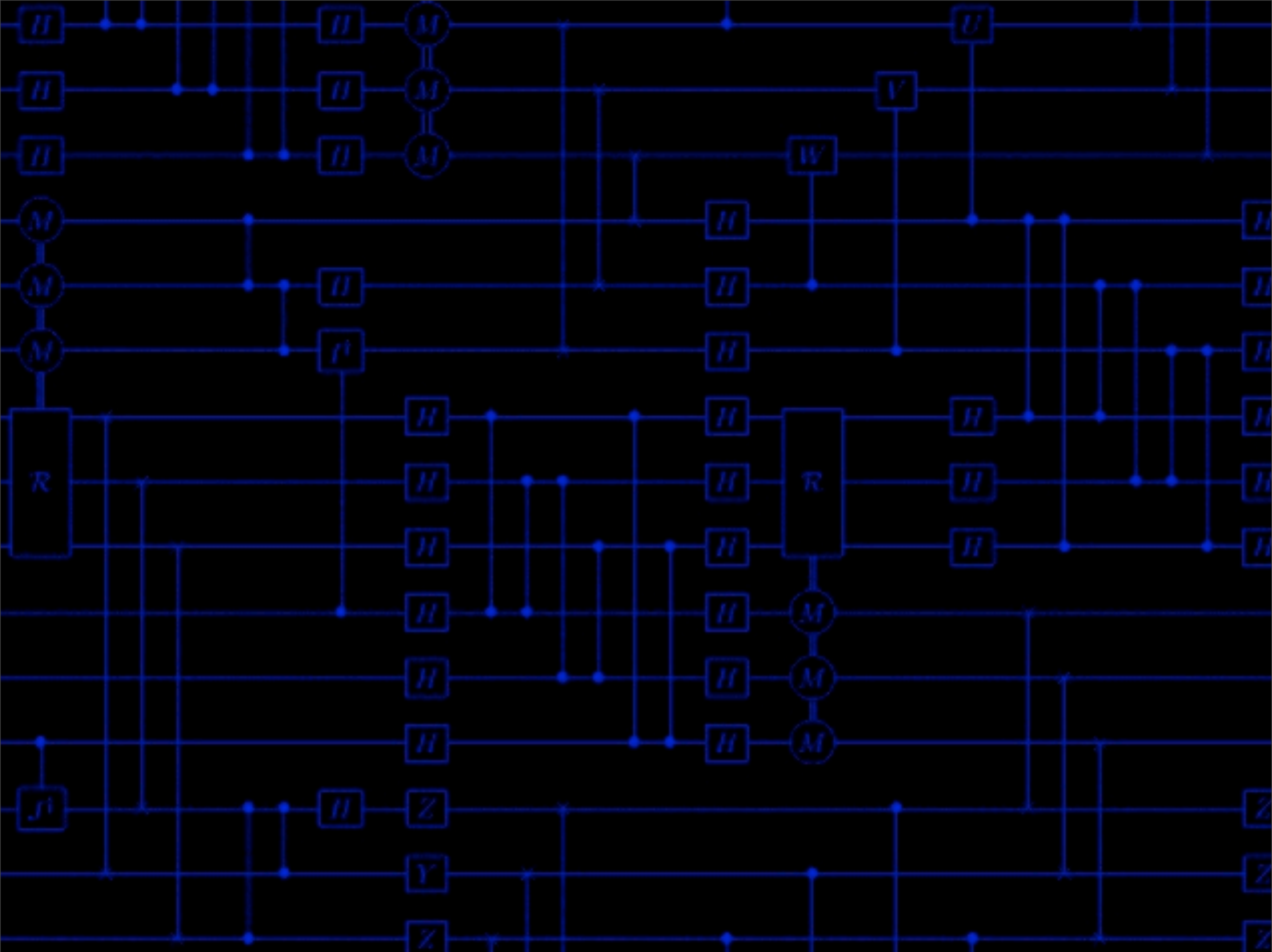
P3. **Reversibility:** Every information processing can be replaced with a reversible one by adding a register in an indivisible initialization.

P4. **Indivisibility of Processing Composition:** The processing corresponding to the input-output sequence of two indivisible processings is itself indivisible.

P5. **Discriminability of Specific Initializations:** For any specific initialization there exists another initialization that can be perfectly discriminated from it.

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What is out of there?



Physics is Information

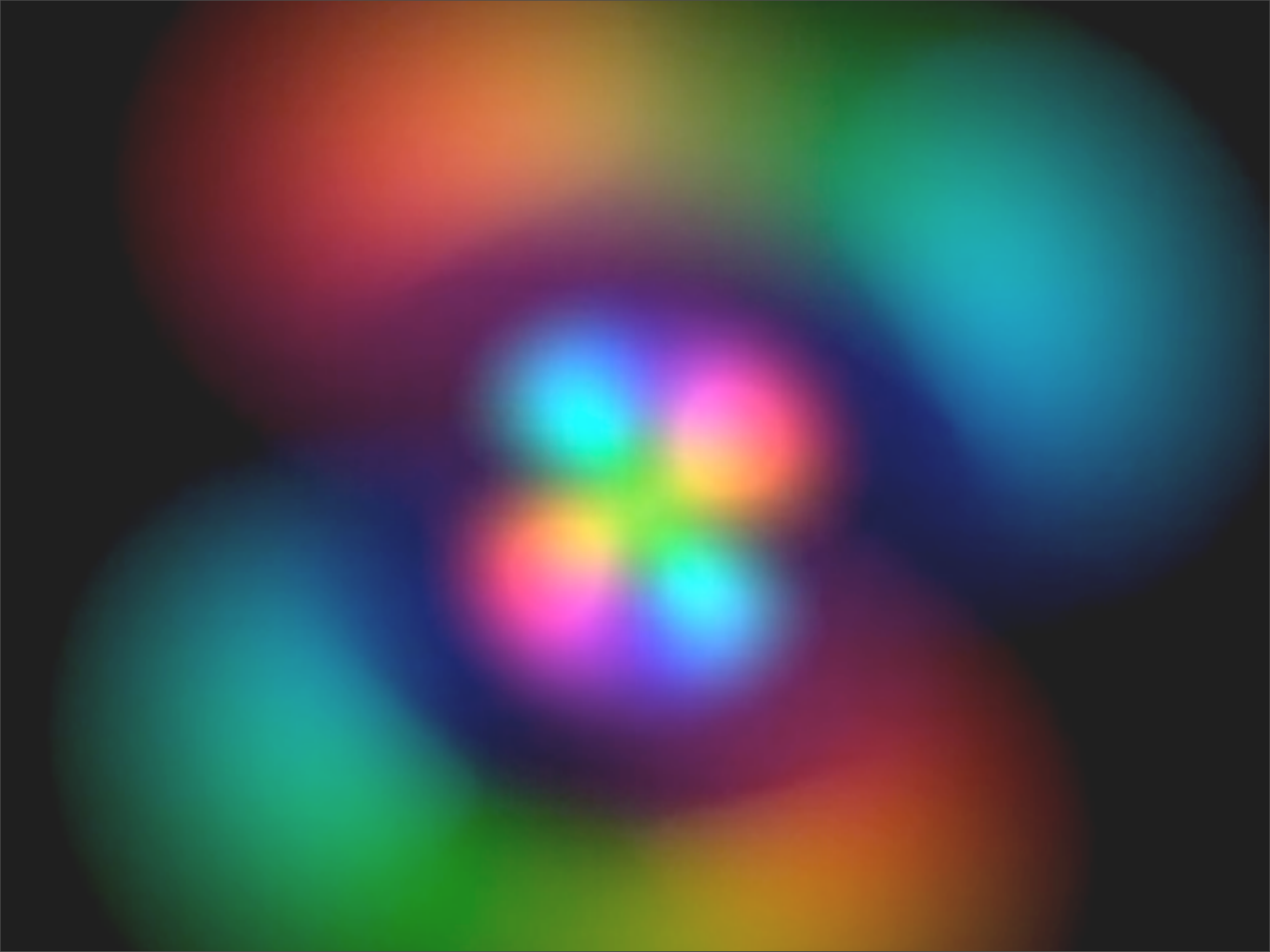
“It from
Bit”



*“Information
is physical”*

(Bit from It)

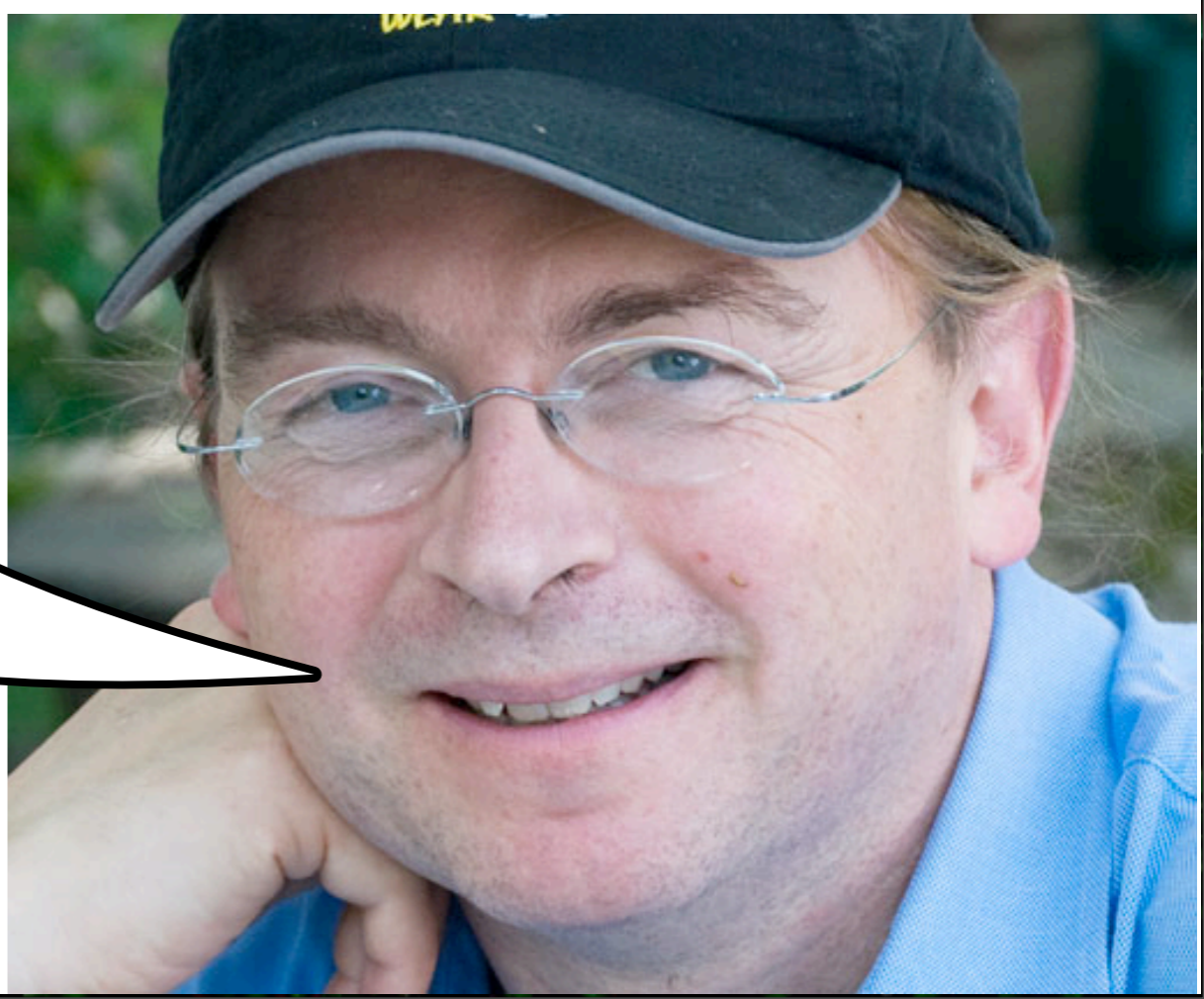




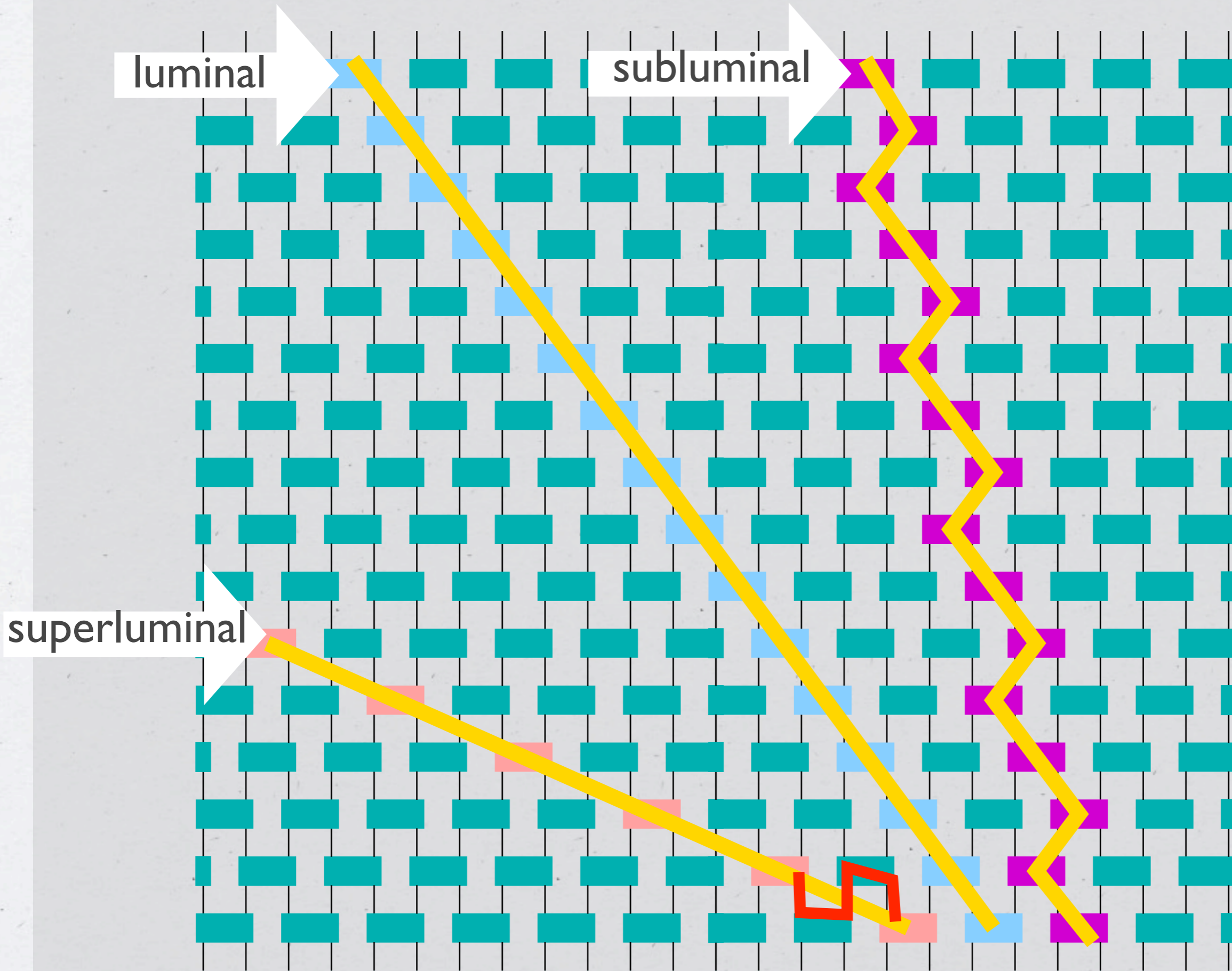
A large table with multiple columns containing alphanumeric data, possibly a ledger or a list of records. The text is oriented vertically on the page.



THE UNIVERSE
IS A HUGE
COMPUTER

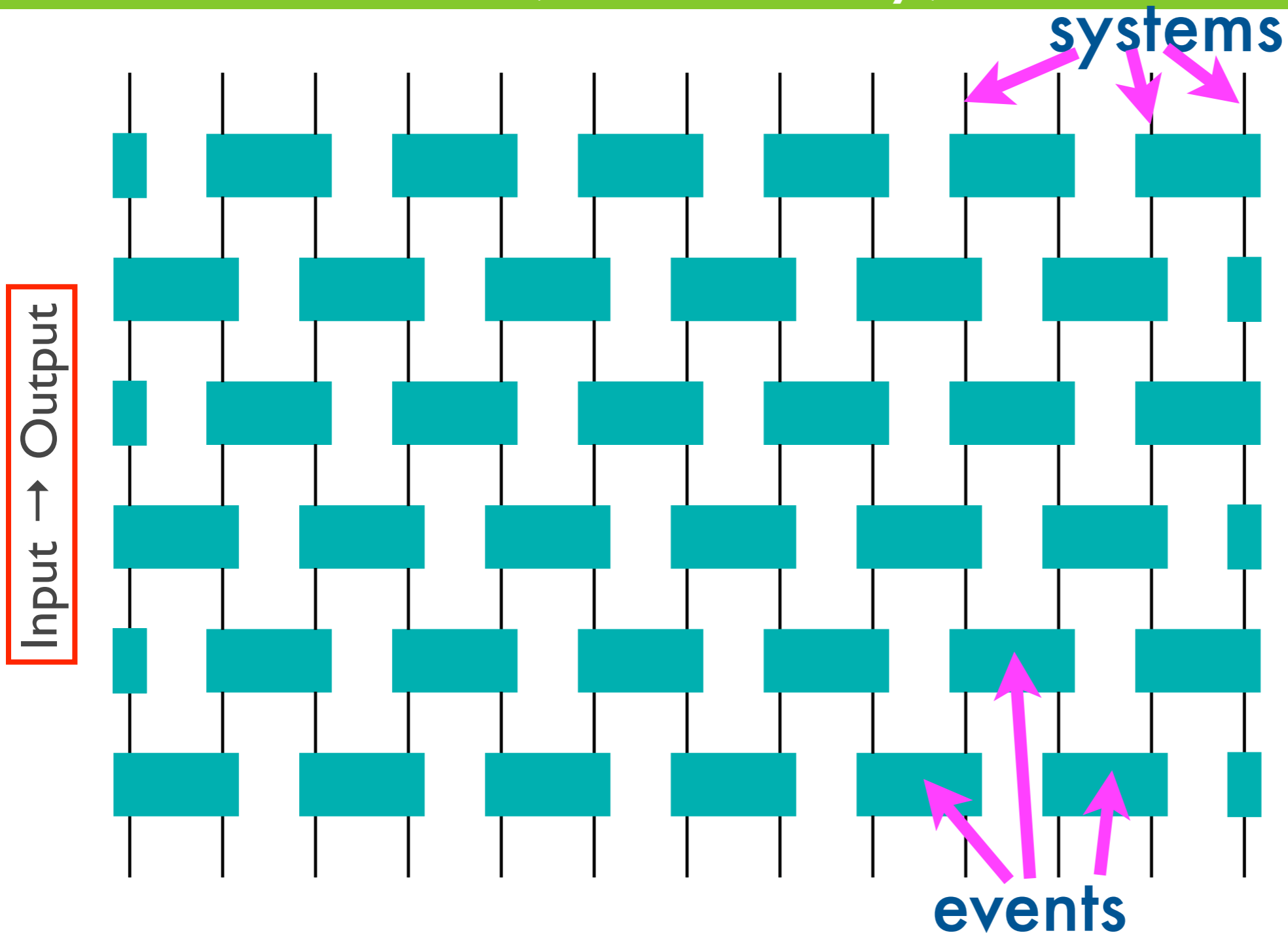


HOW RELATIVITY EMERGES FROM THE COMPUTATION?



Relativity from QT

(from causality)

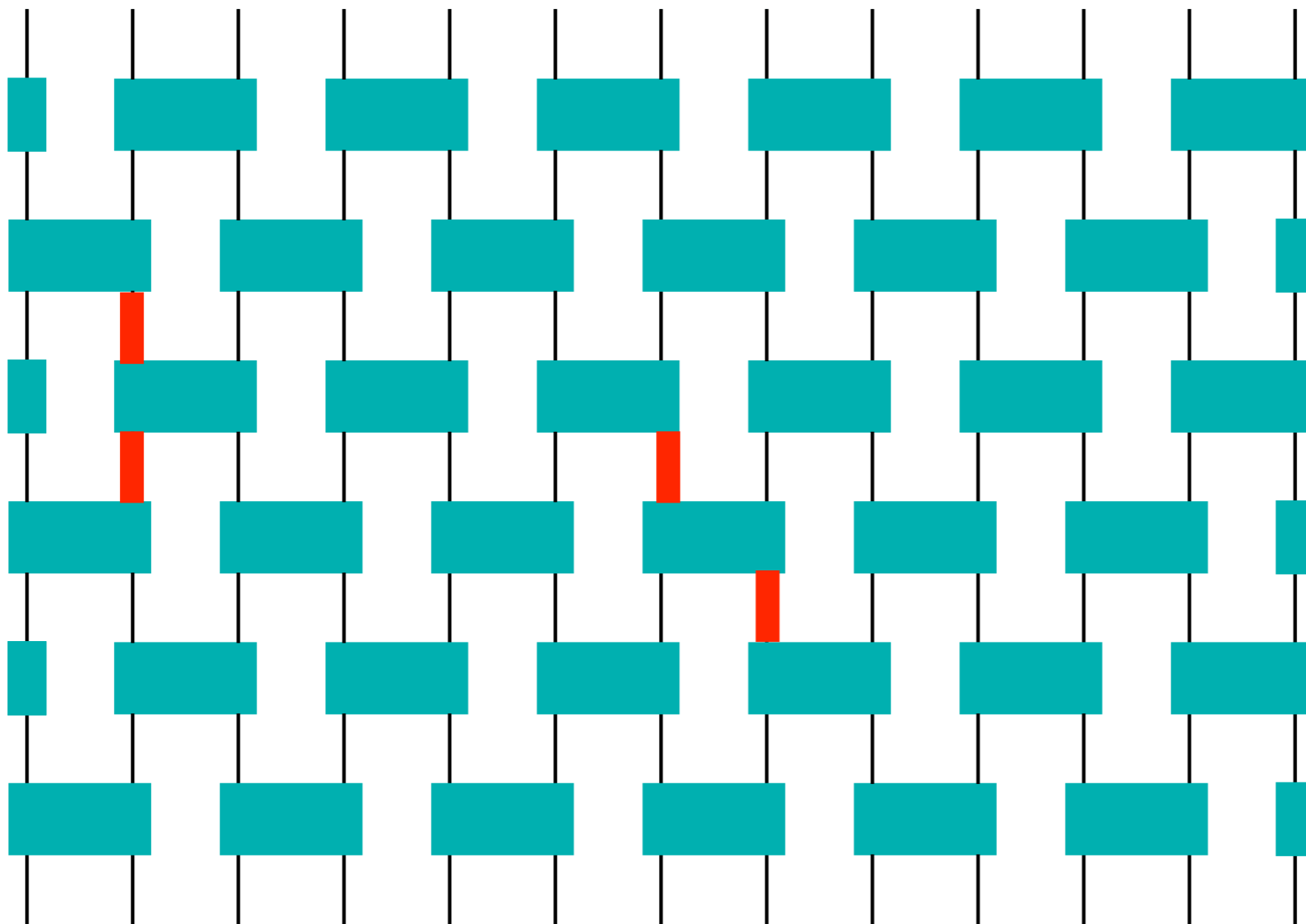


Relativity from QT

(from causality)

causally
connected
systems

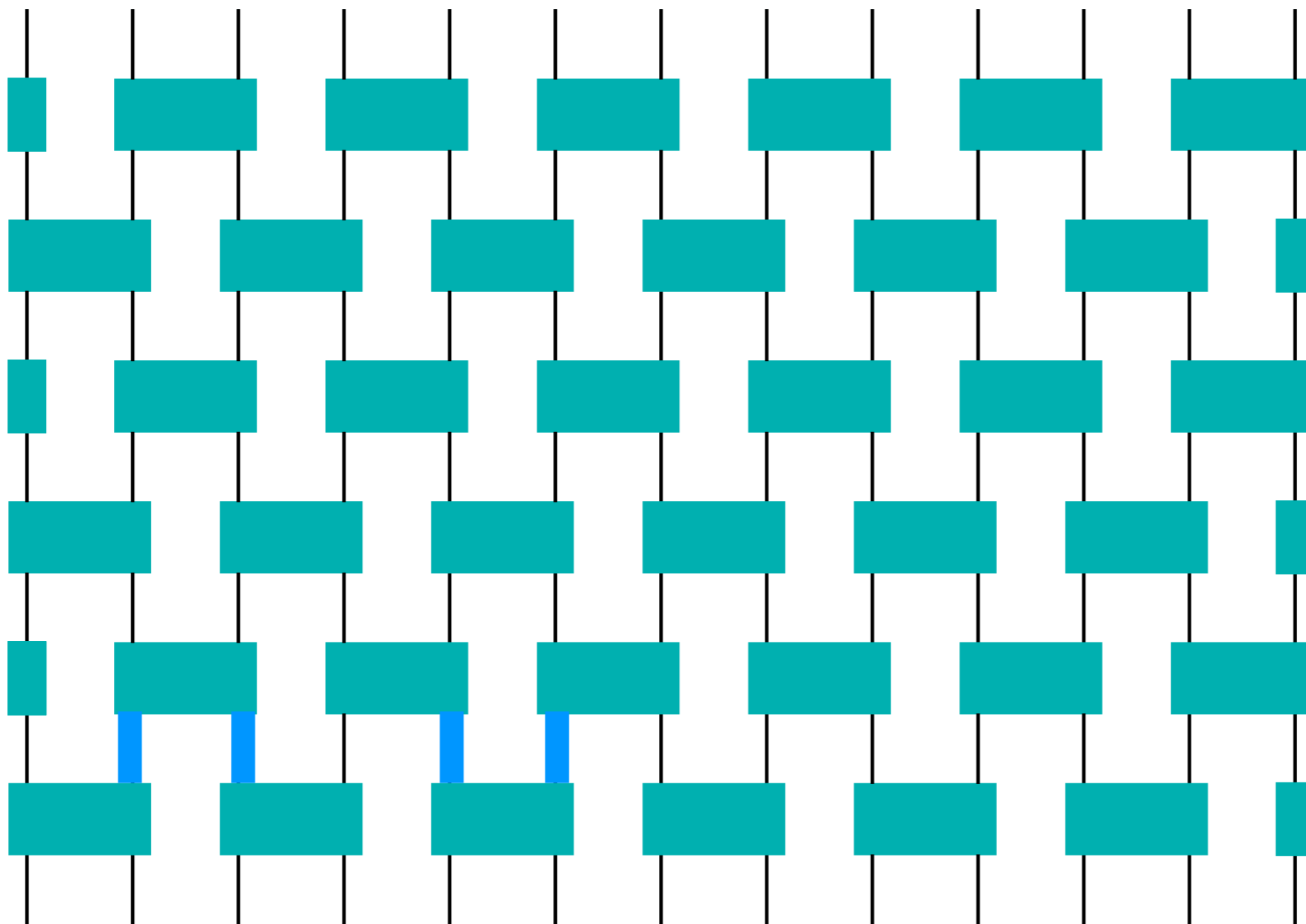
Input → Output



Relativity from QT

(from causality)

Input → Output



independent
systems

“slice”

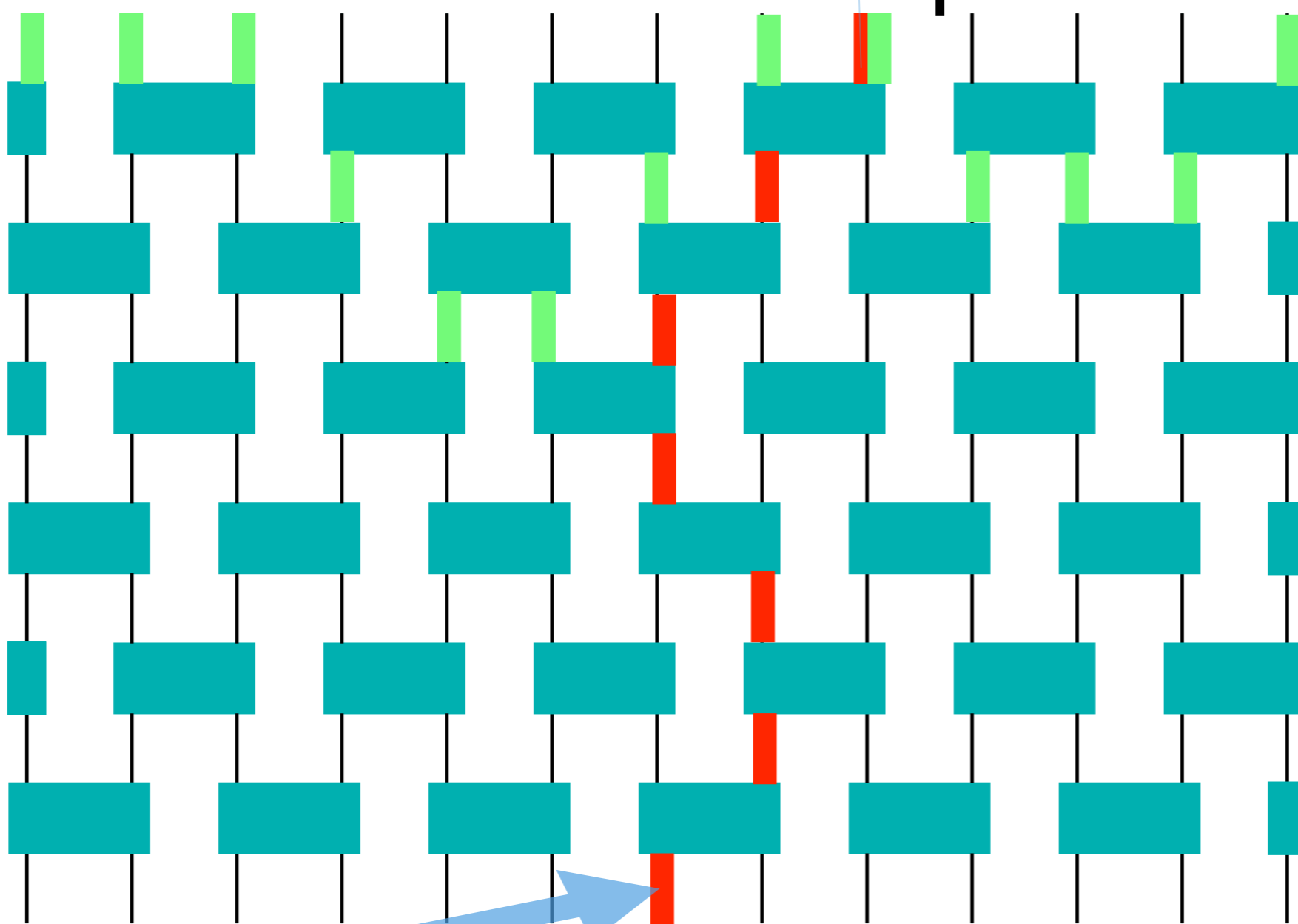
Relativity from QT

(from causality)

topology
(Alexandrov)
metric =
event-counting

causal antichain = space

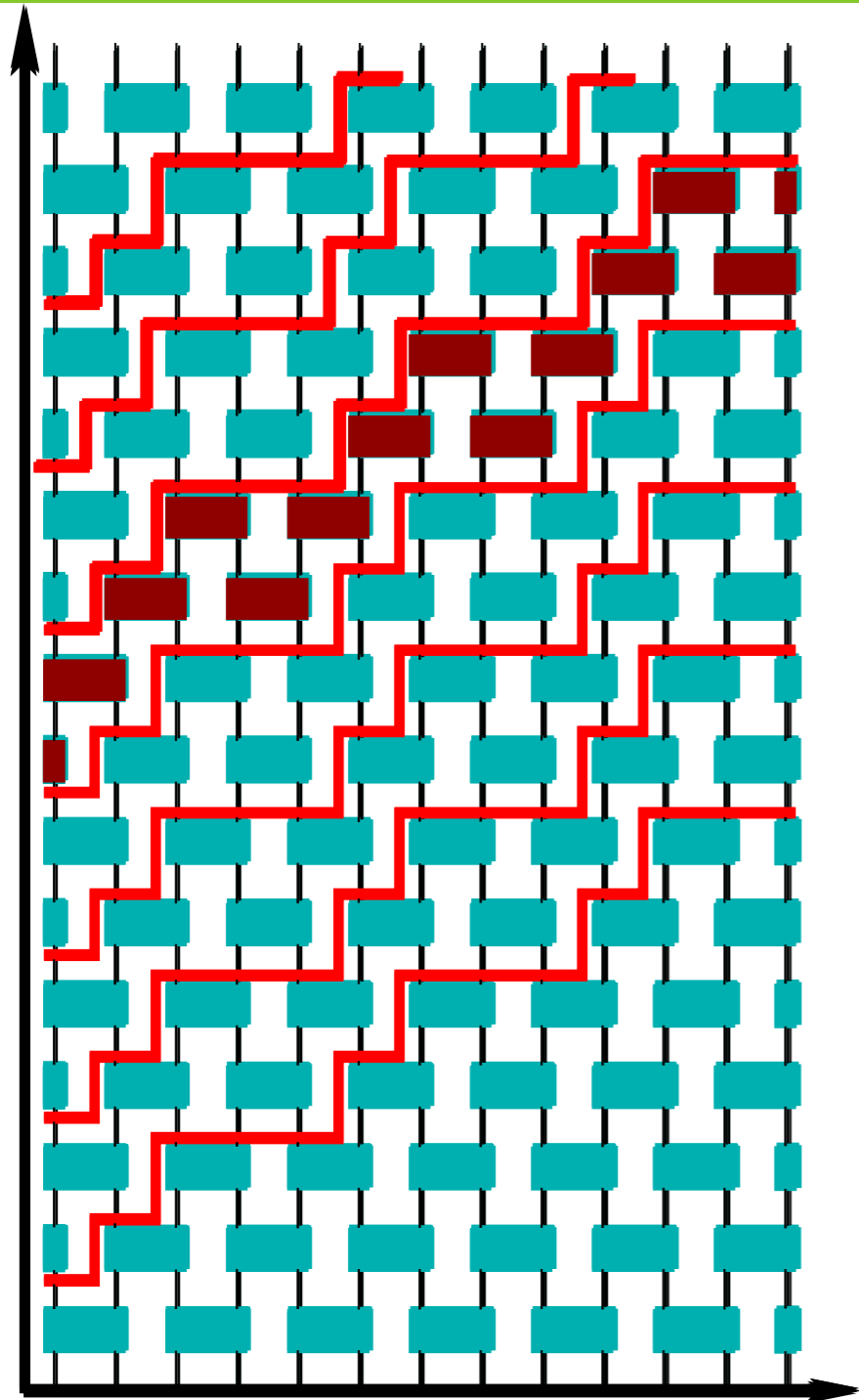
Input → Output



causal chain = time (observer)

Relativity from QT

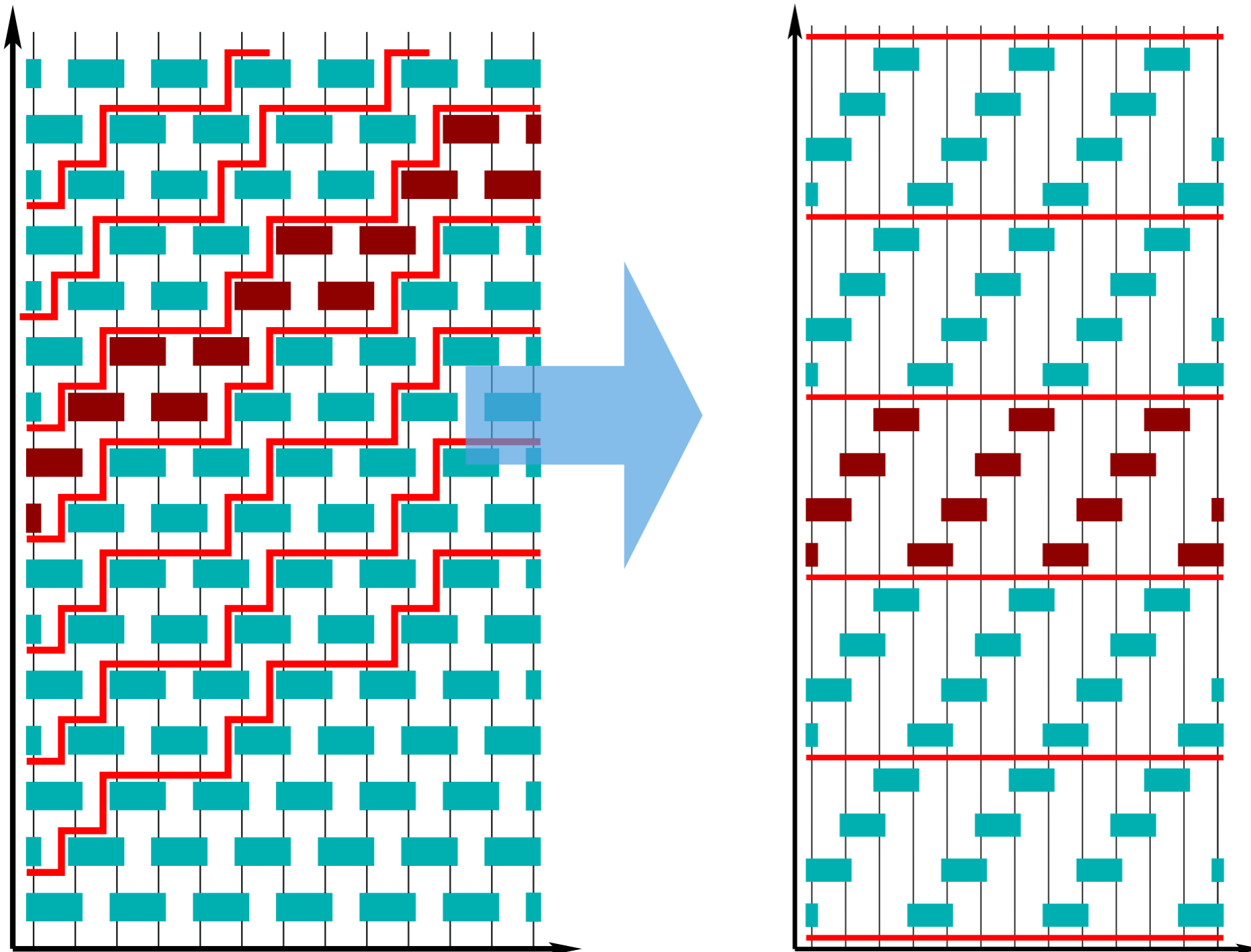
(from causality)



build a
uniform
foliation

Relativity from QT

(from causality)

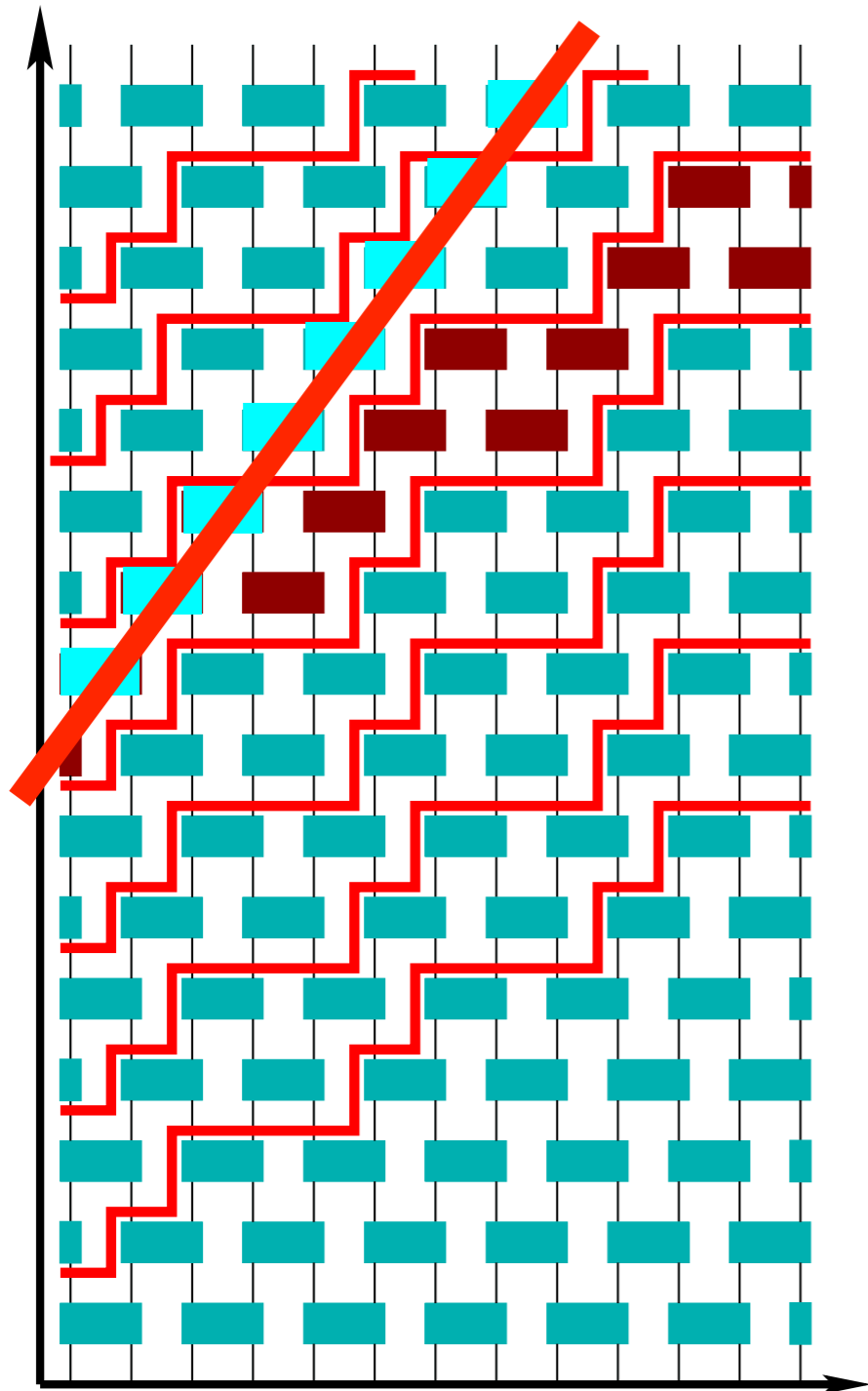


change
reference

Relativity from QT

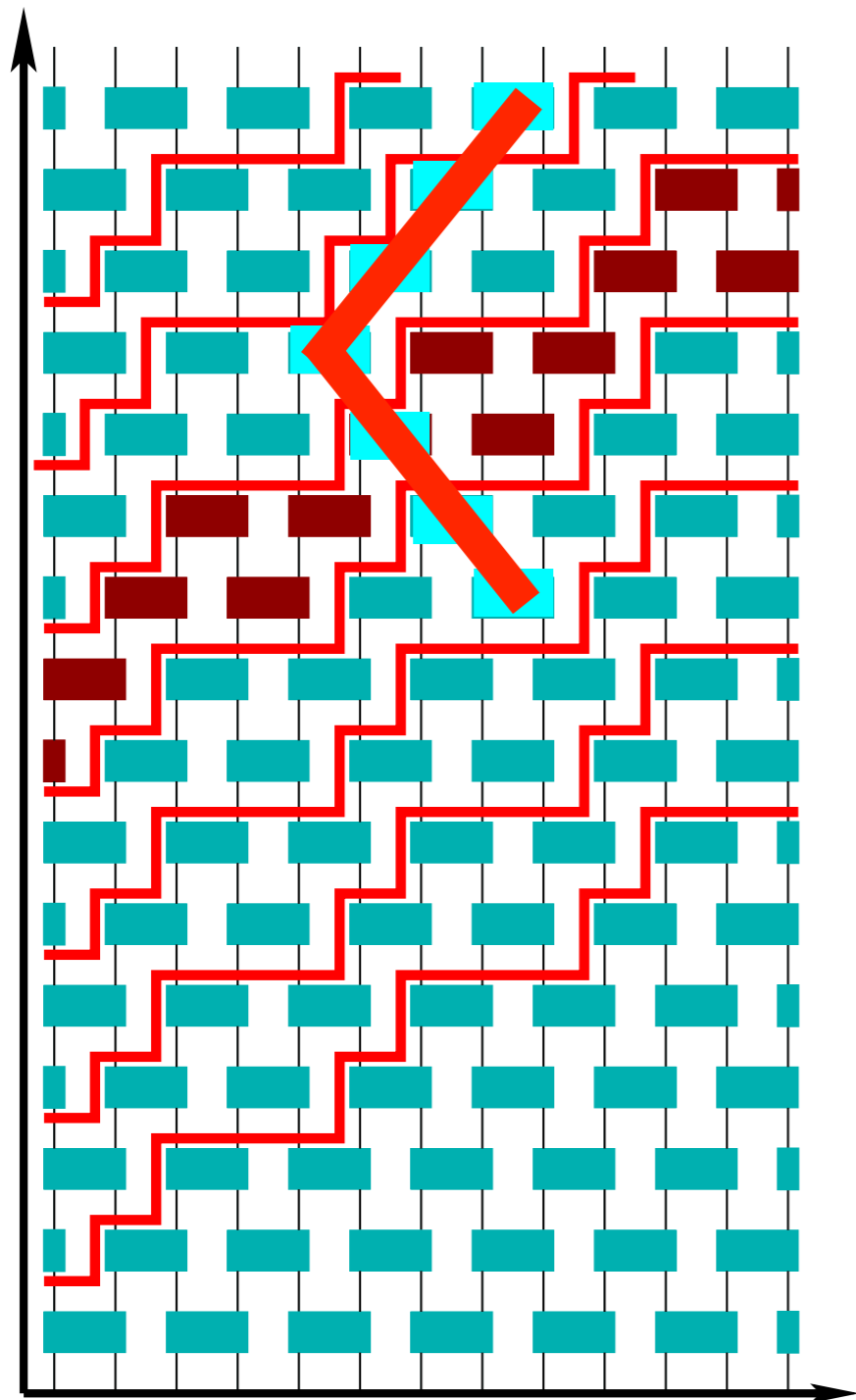
(from causality)

speed of
light



Relativity from QT

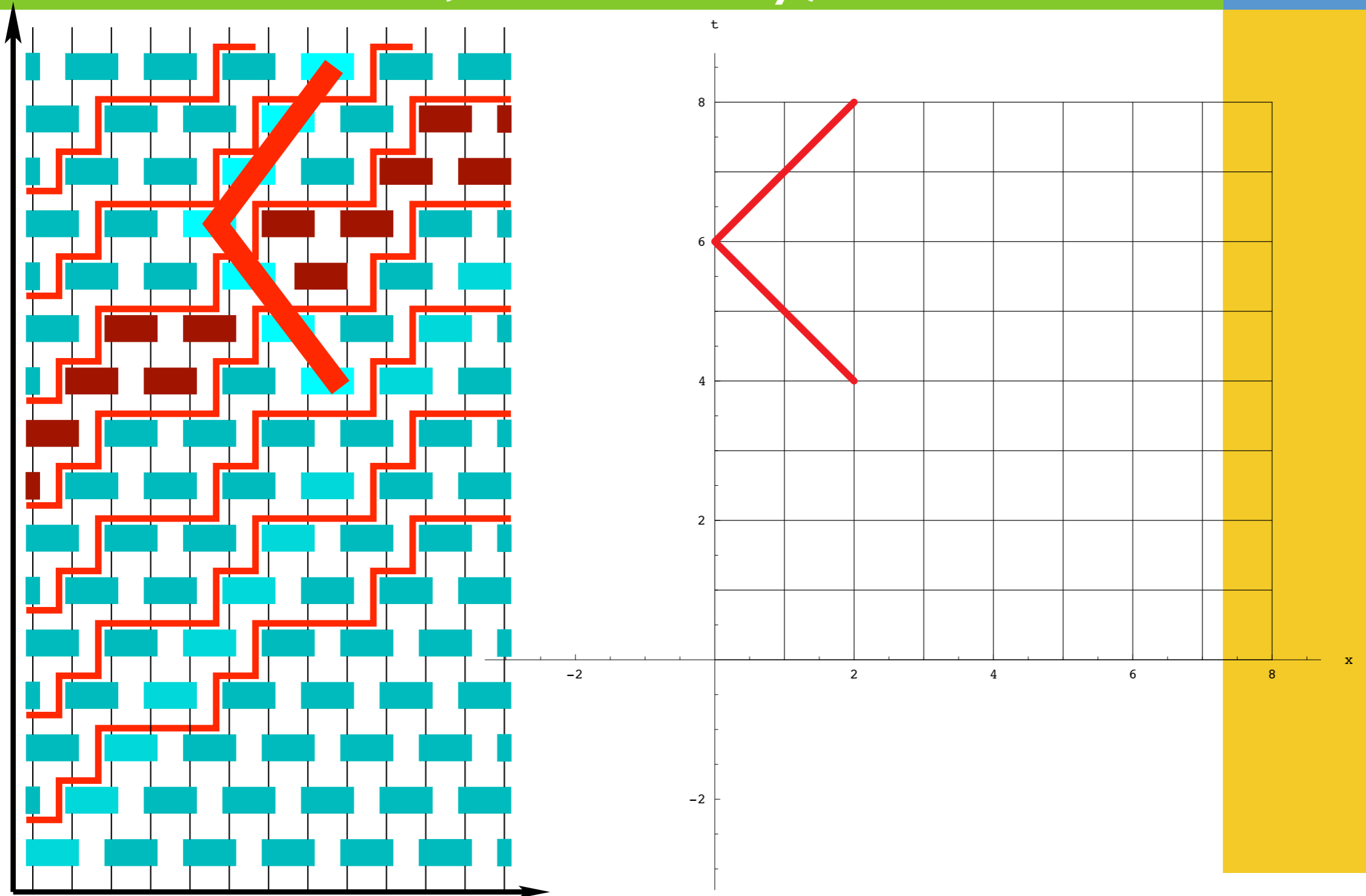
(from causality)



clock tic-tac

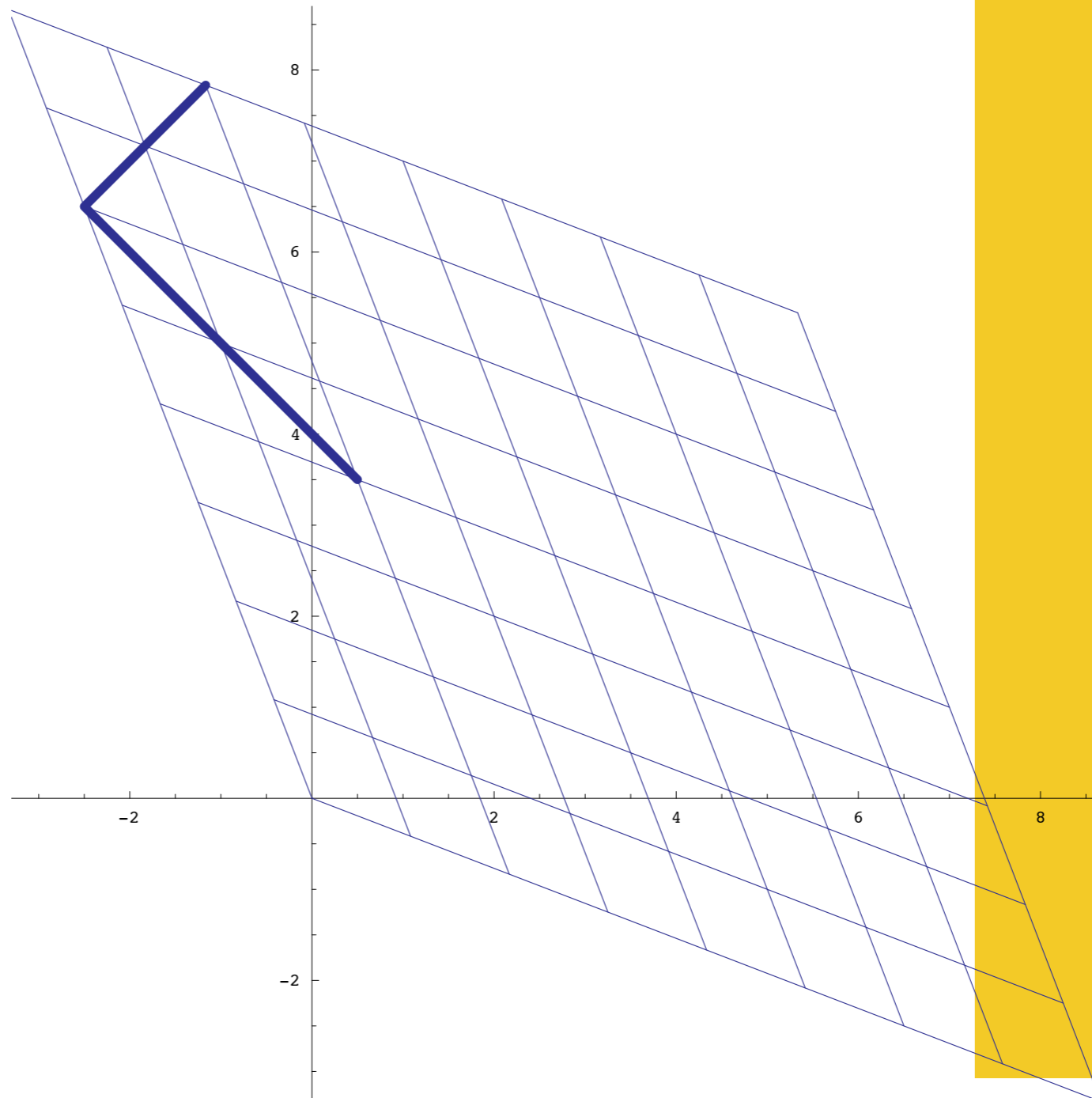
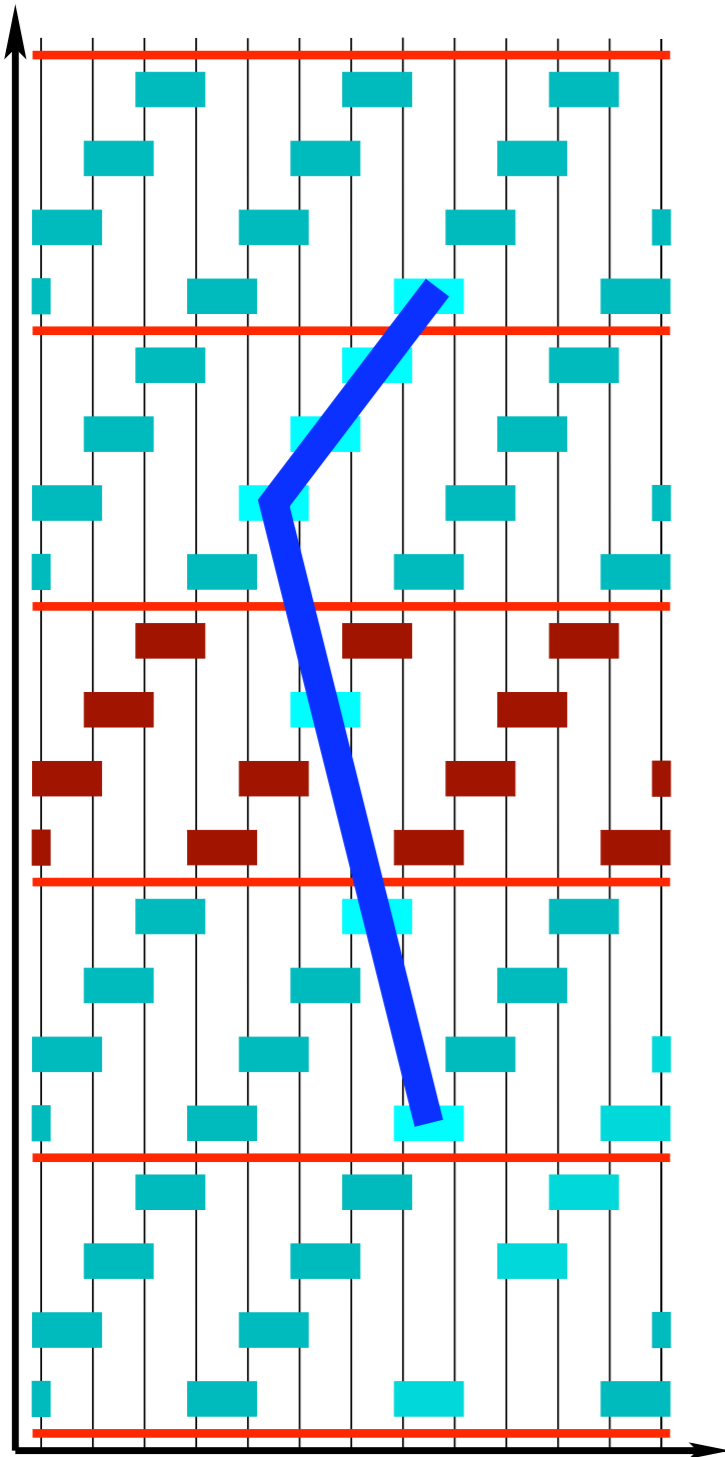
Relativity from QT

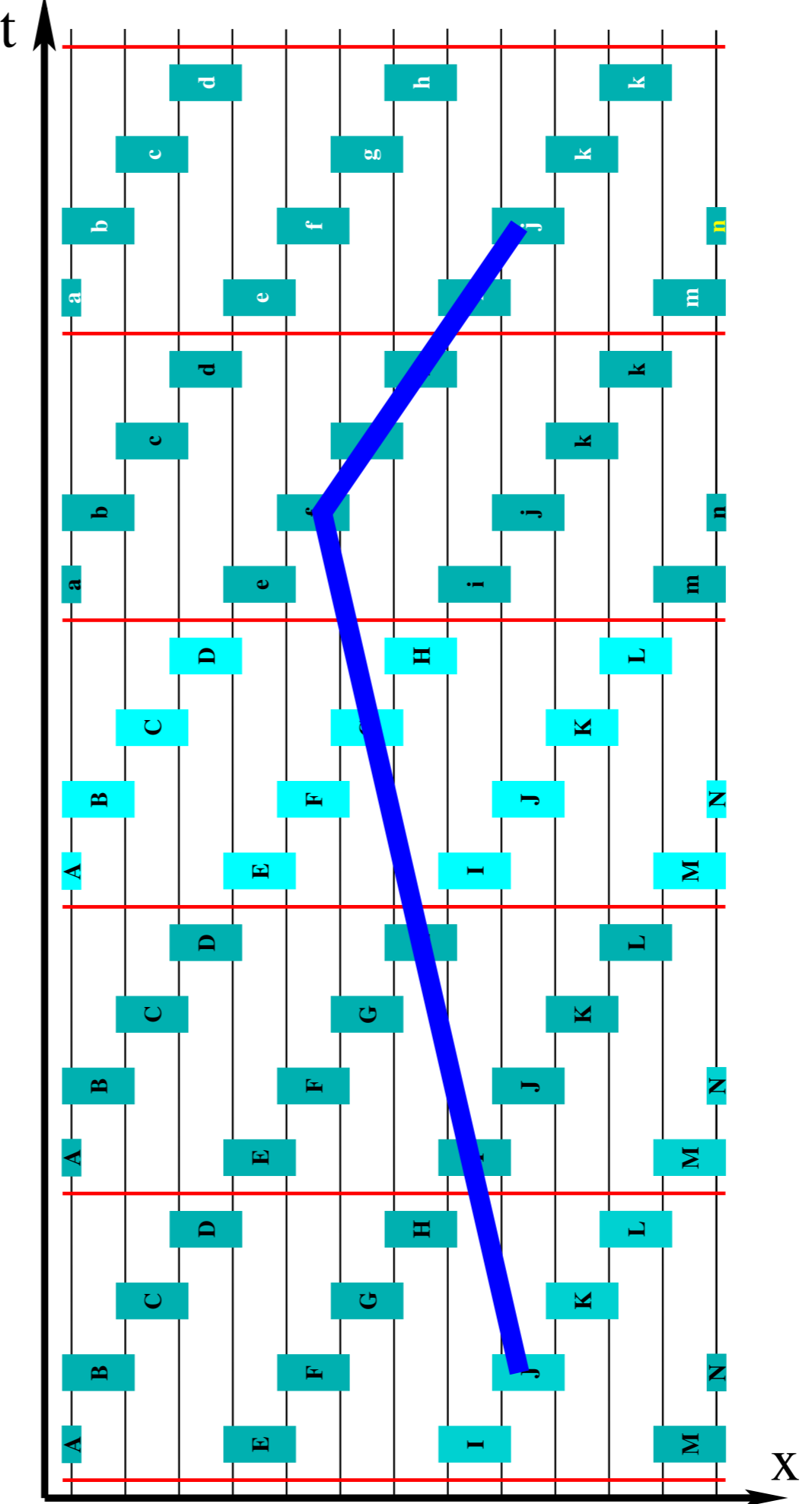
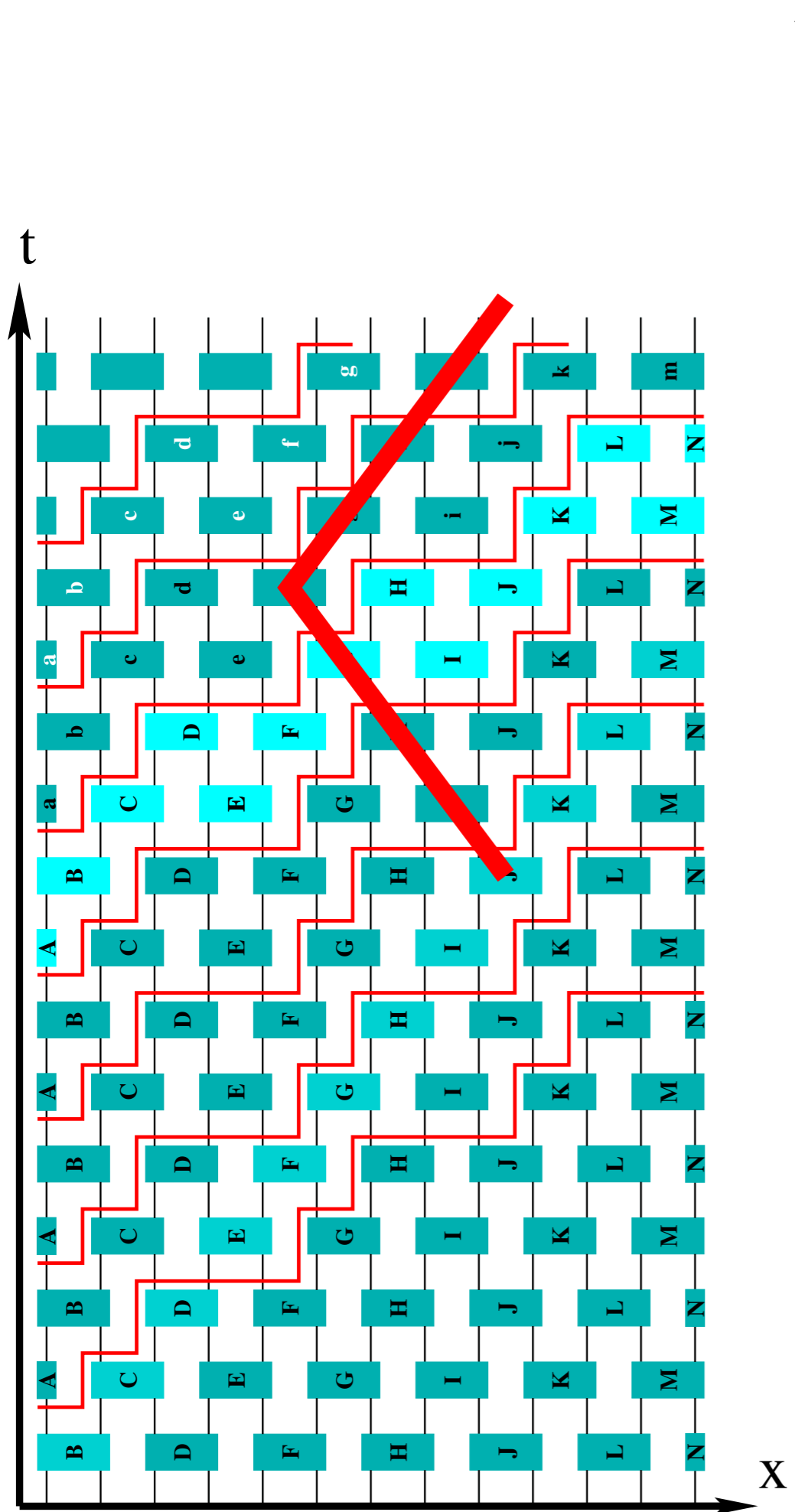
(from causality)



Relativity from QT

(from causality)

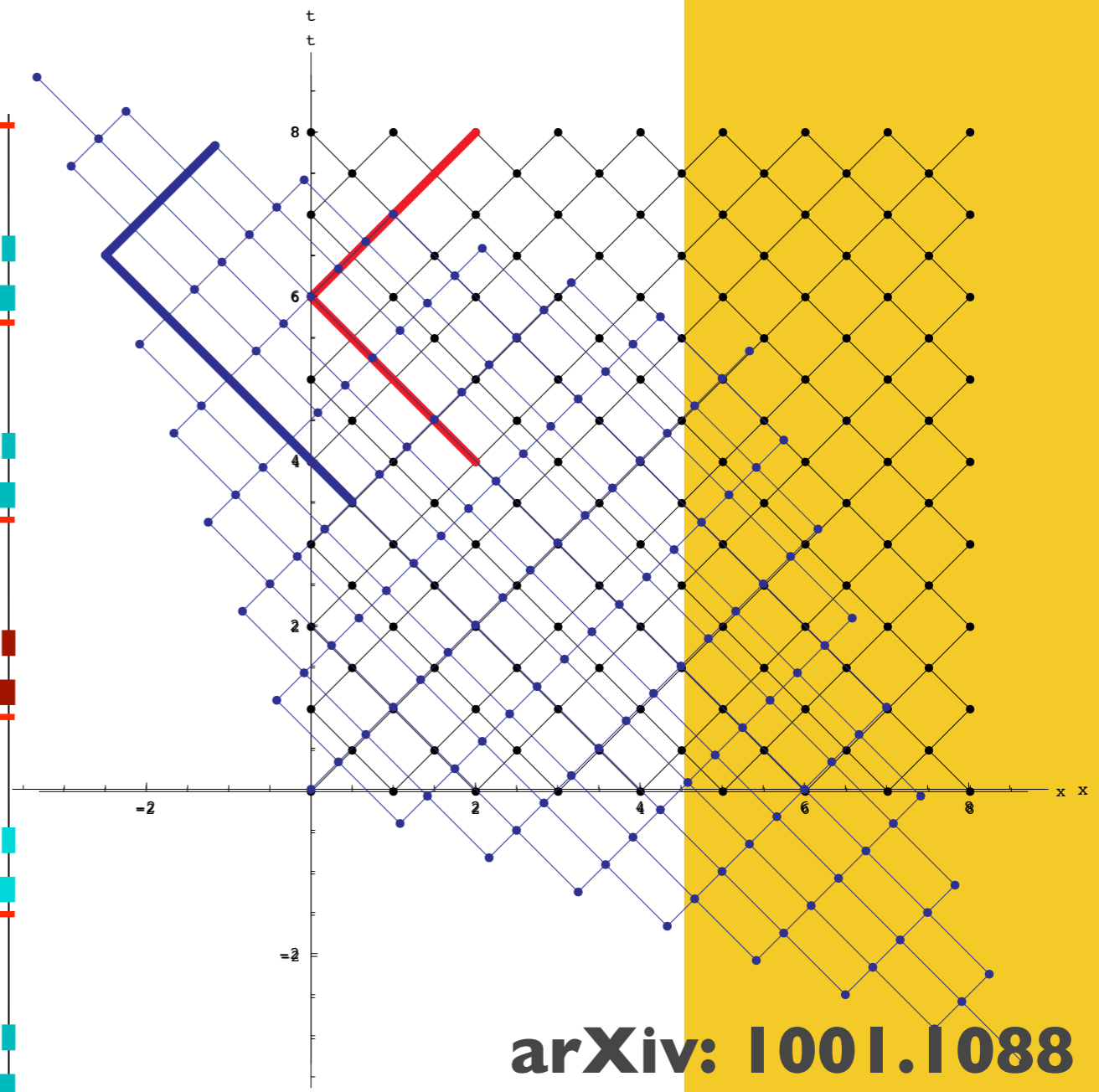
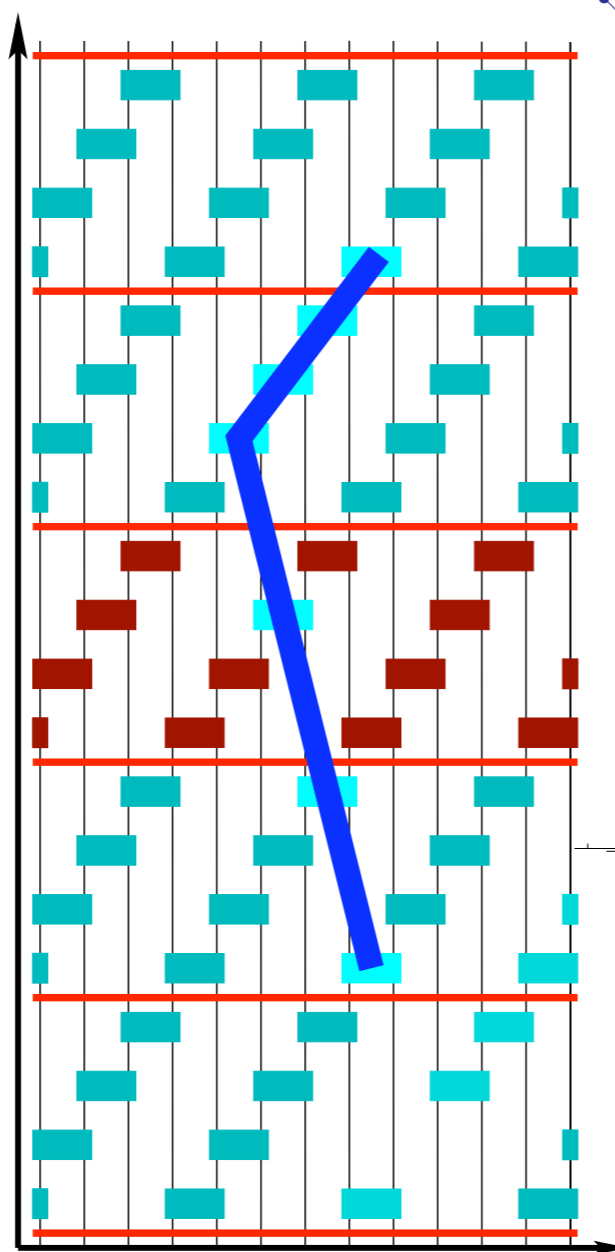
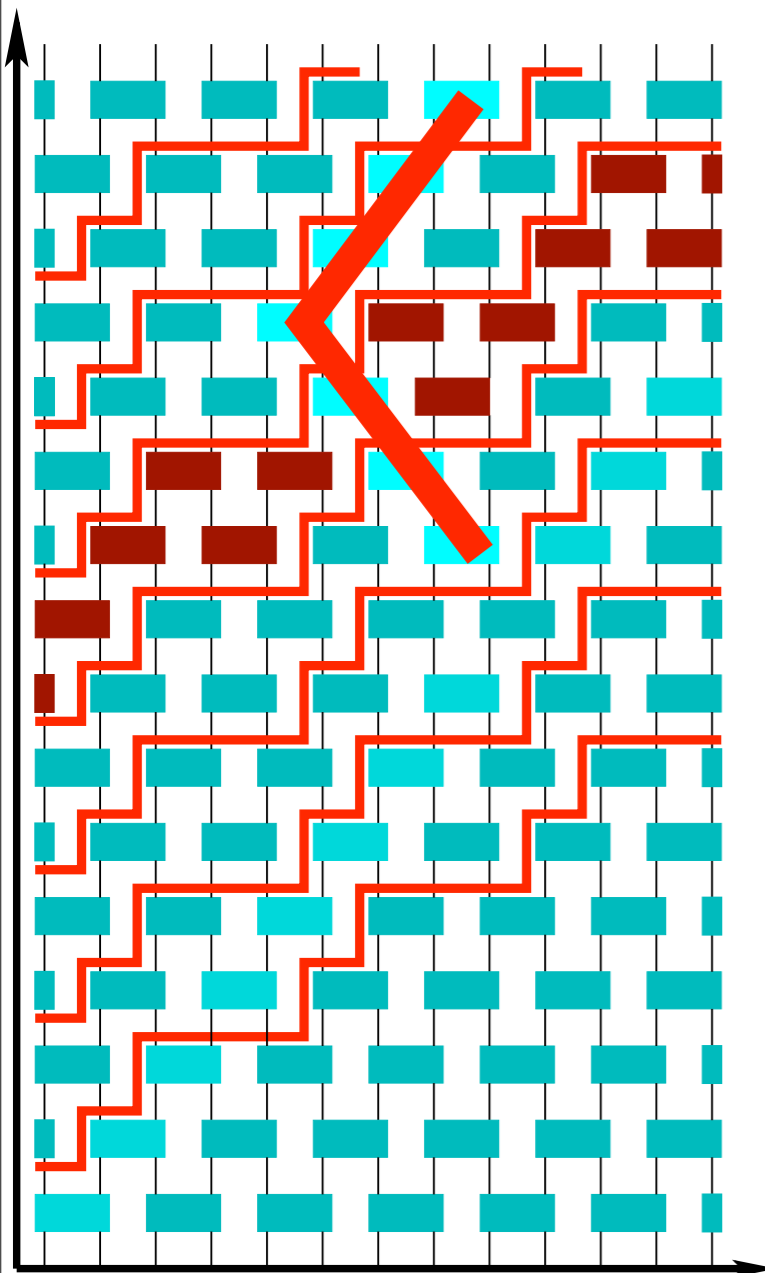




Time dilation and space contraction

Relativity from QT

(from causality)



arXiv: 1001.1088

WE GOT RELATIVITY FROM
CAUSALITY!

WE GOT MUCH MORE:

FROM CAUSALITY WE GOT
SPACE AND TIME ENDOWED
WITH RELATIVITY!

Relativity from QT

A theory of quantum gravity based on quantum computation

Seth Lloyd

Massachusetts Institute of Technology

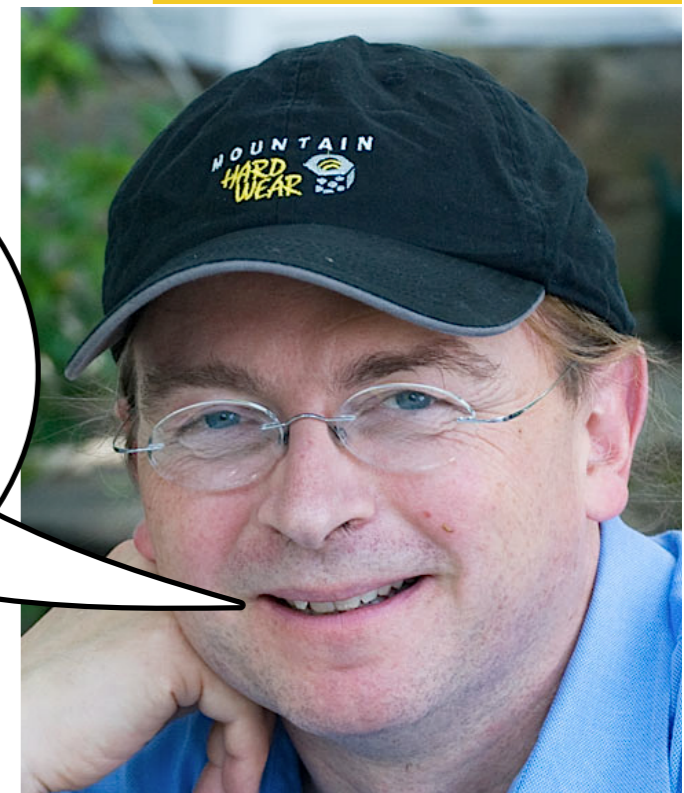
MIT 3-160, Cambridge, Mass. 02139 USA

slloyd@mit.edu

Keywords: quantum computation, quantum gravity

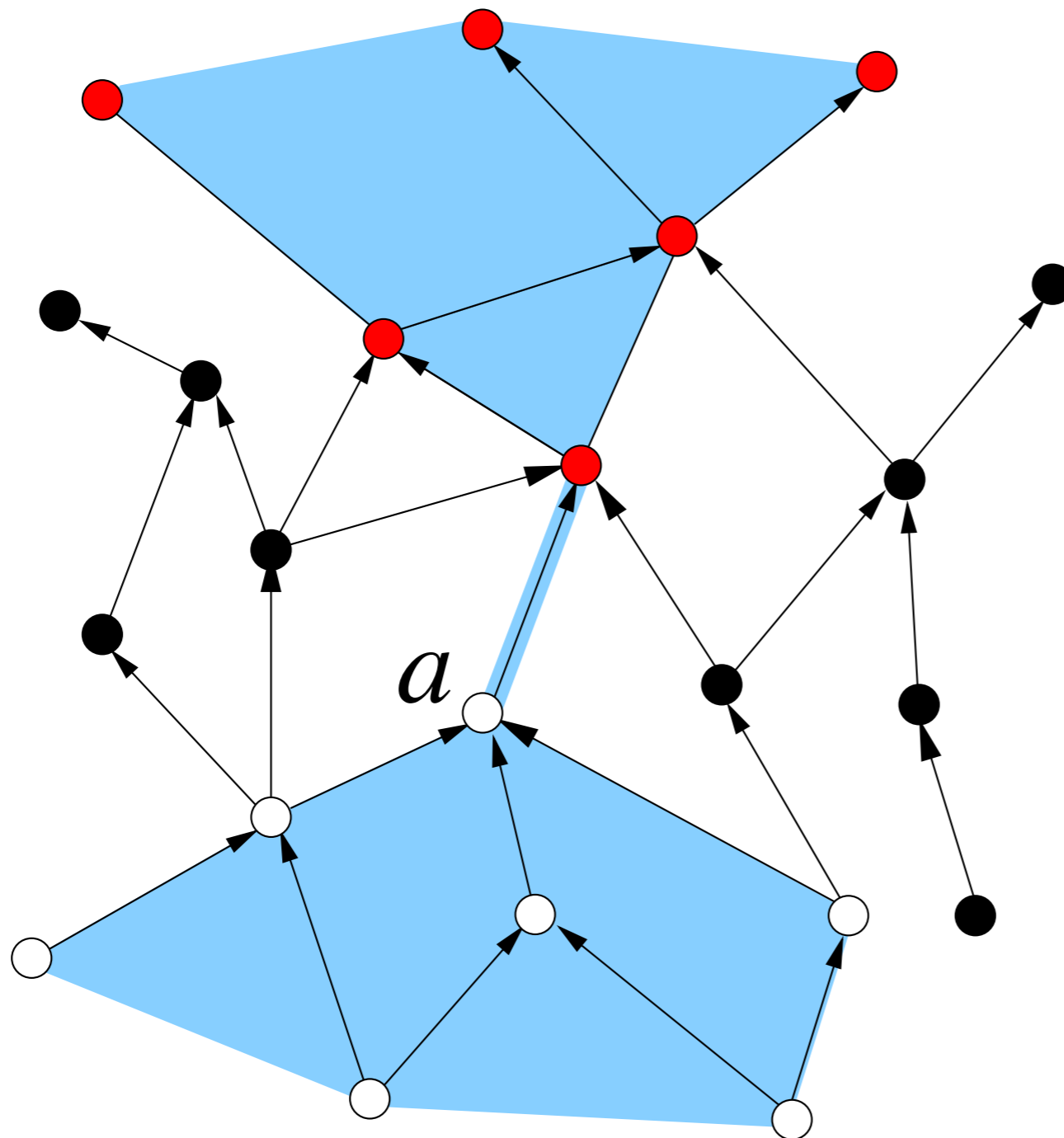
Abstract: This paper proposes a method of unifying quantum gravity with quantum computation. In this theory, fundamental spacetime geometry is a construct derived from the underlying quantum computation. Pairwise interactions between quantum degrees of freedom give rise to a superposition of four-geometries that obeys the Einstein-Regge equations. The theory predicts a reaction of the metric to computational 'matter,' based on quantum cosmology.

THE GEOMETRY OF
SPACE-TIME IS A
CONSTRUCT DERIVED
FROM THE UNDERLYING
QUANTUM INFORMATION
PROCESSING



Lorentz transformations from causality and topological homogeneity

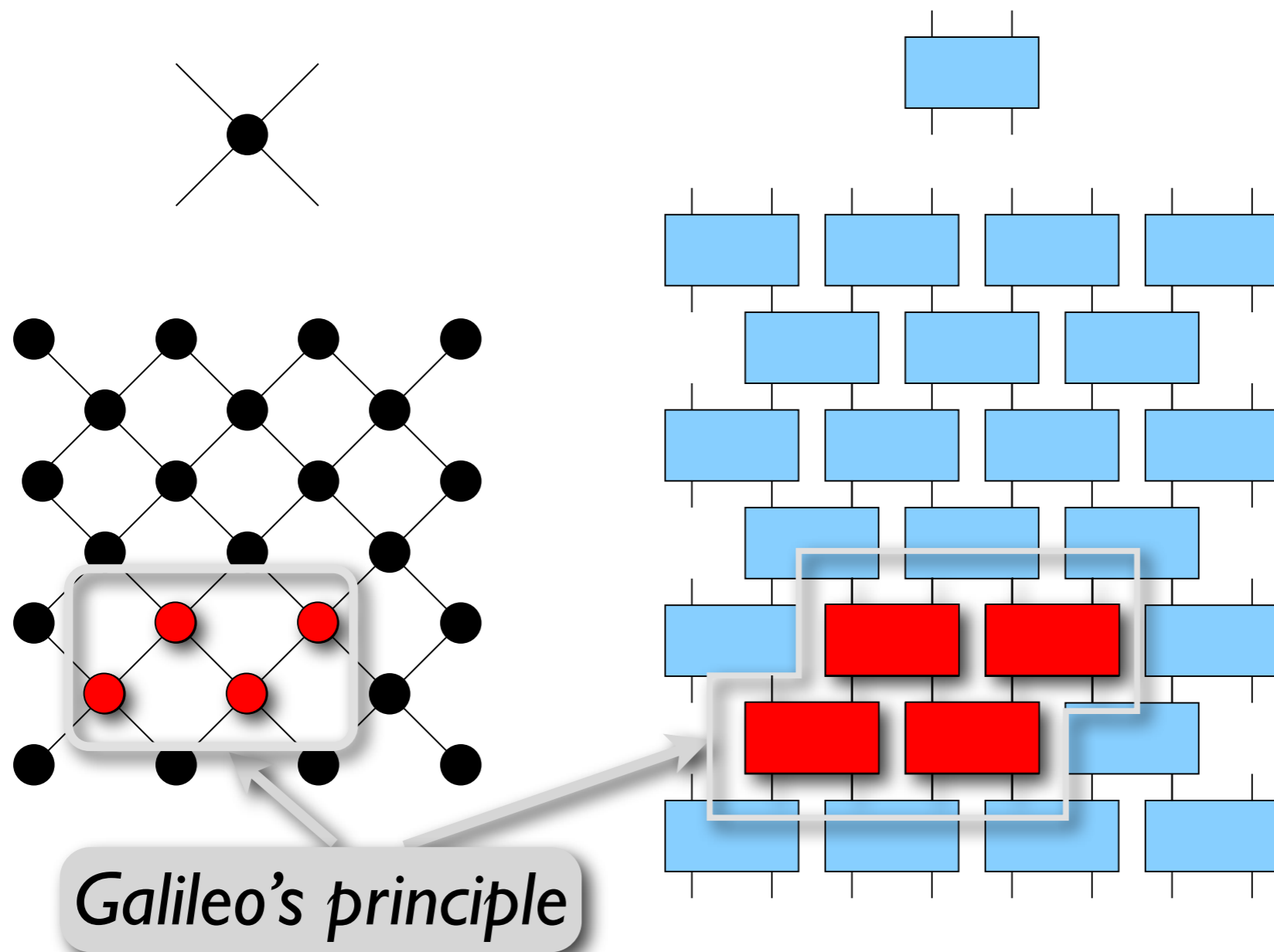
GMD and
A. Tosini
1008.4805



“Light” cones

Lorentz transformations from causality and topological homogeneity

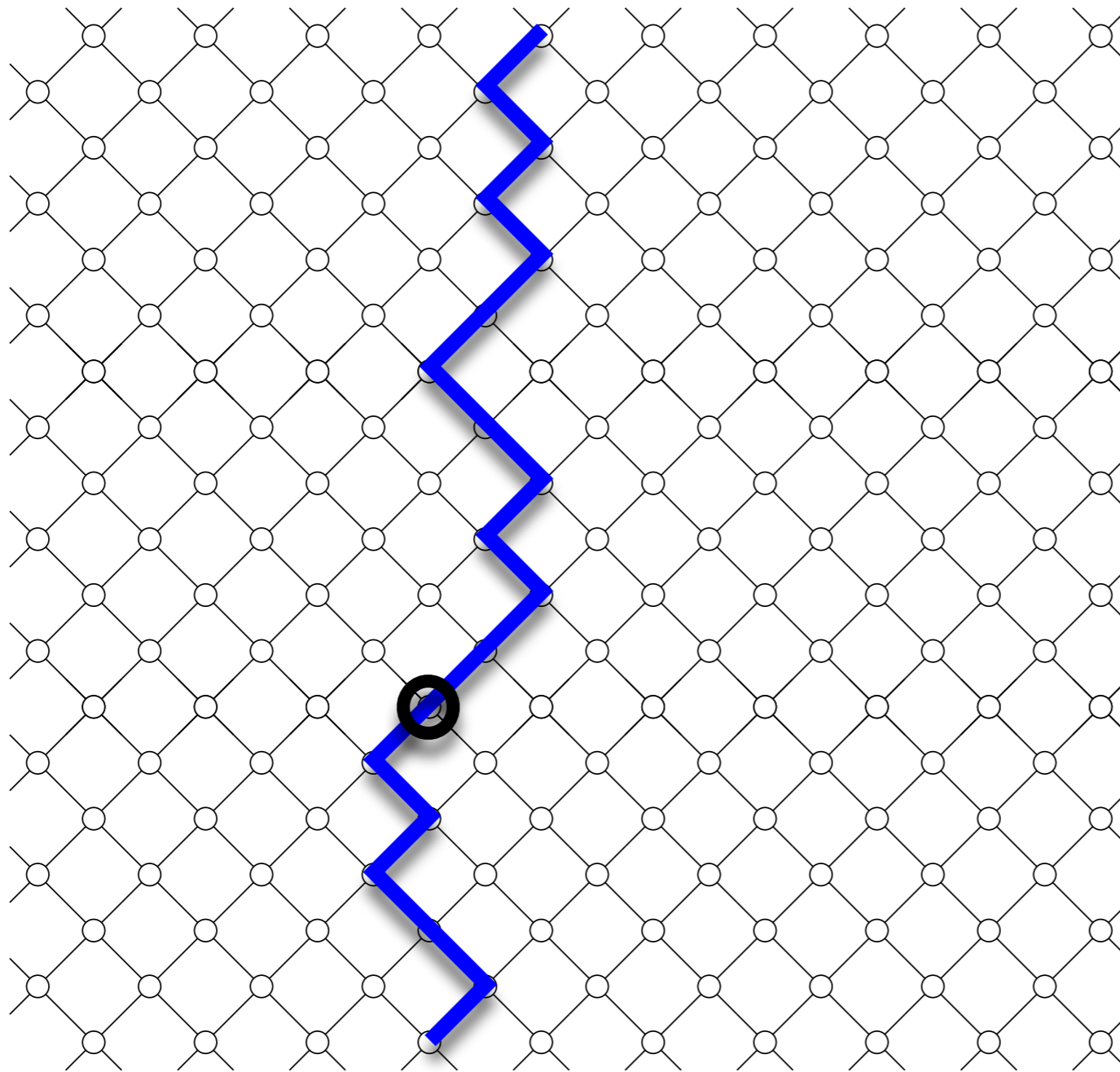
GMD and
A. Tosini
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topological homogeneity

Lorentz transformations from causality and topological homogeneity

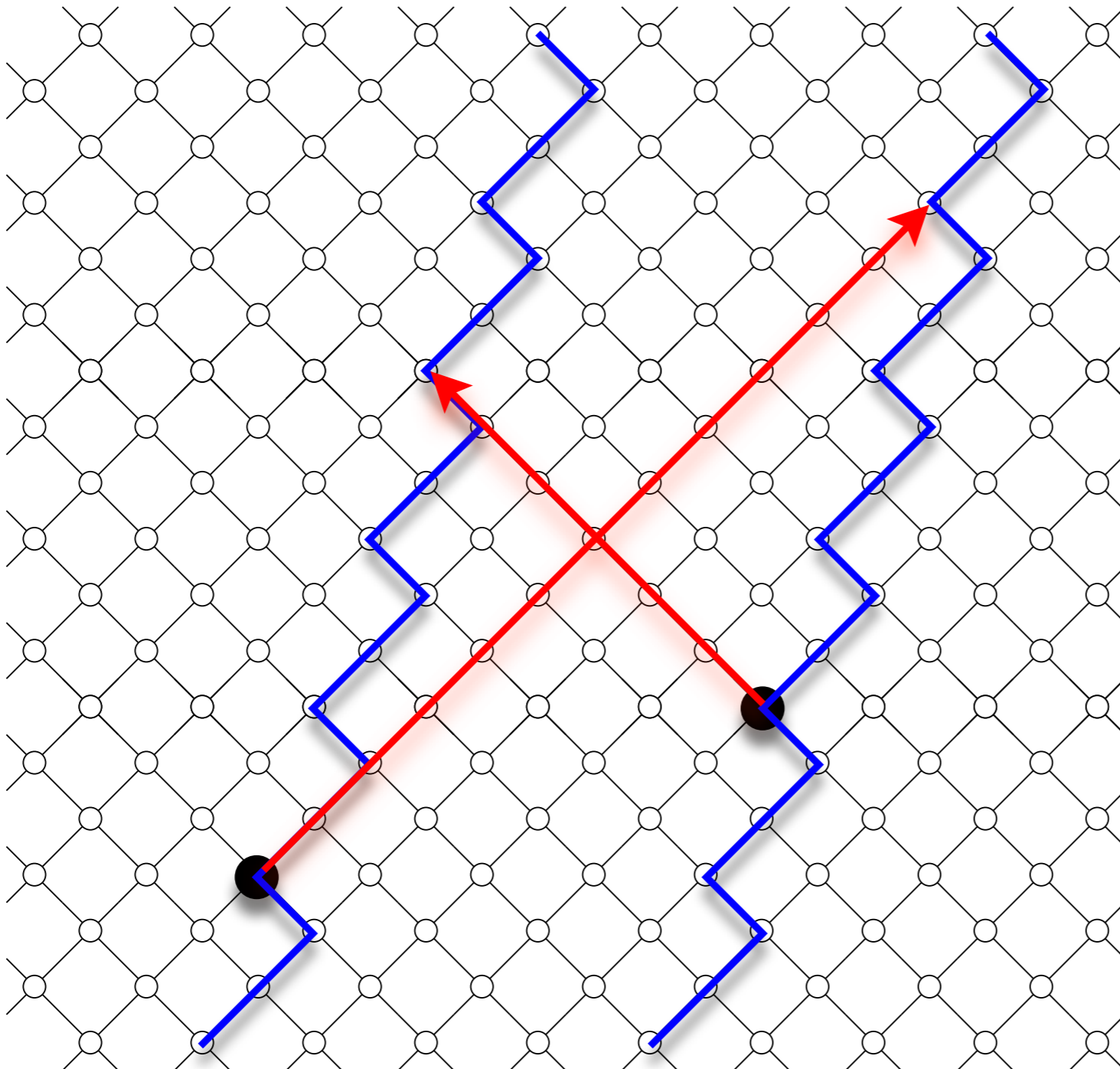
GMD and
A. Tosini
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observer

Lorentz transformations from causality and topological homogeneity

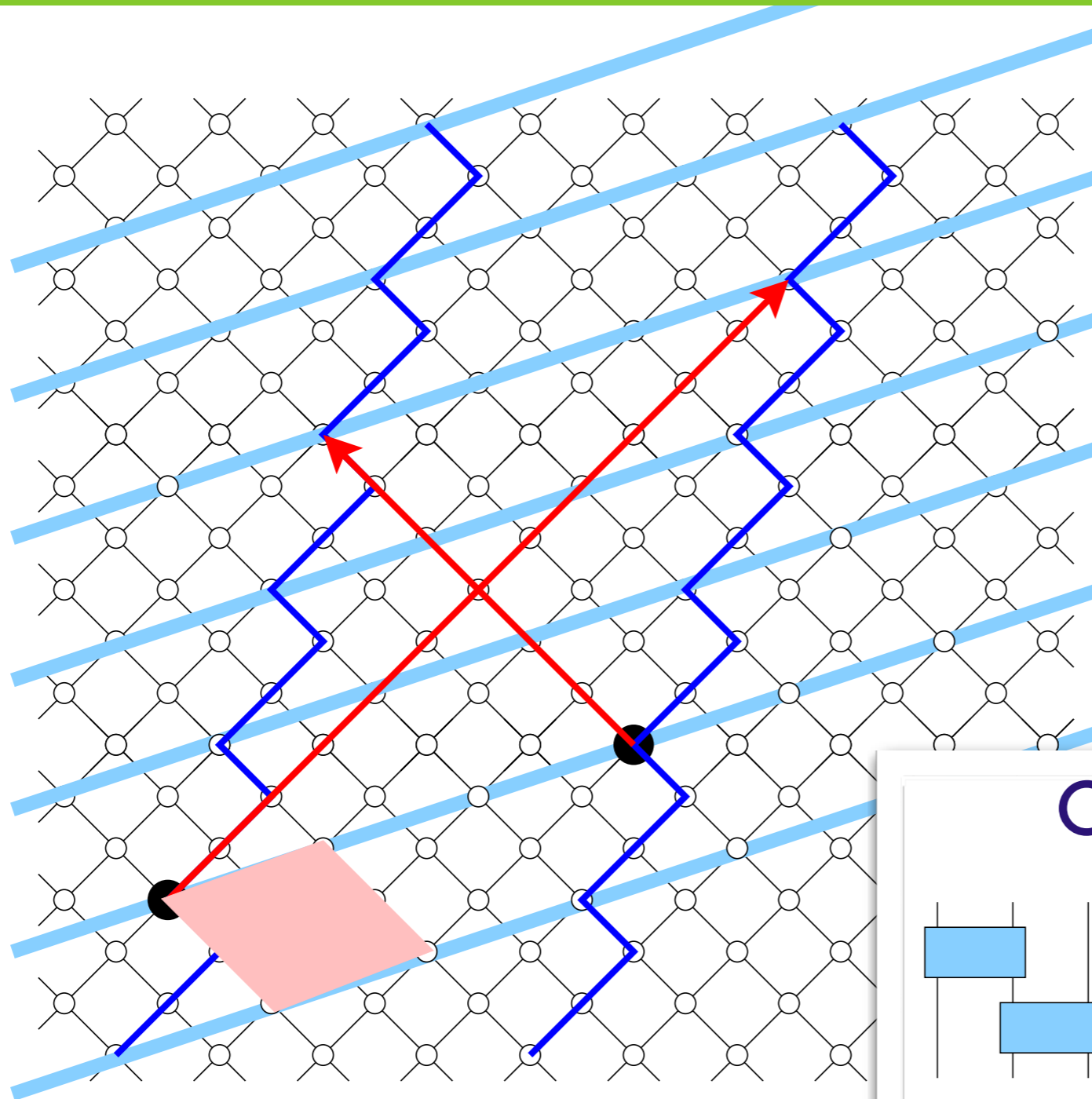
GMD and
A. Tosini
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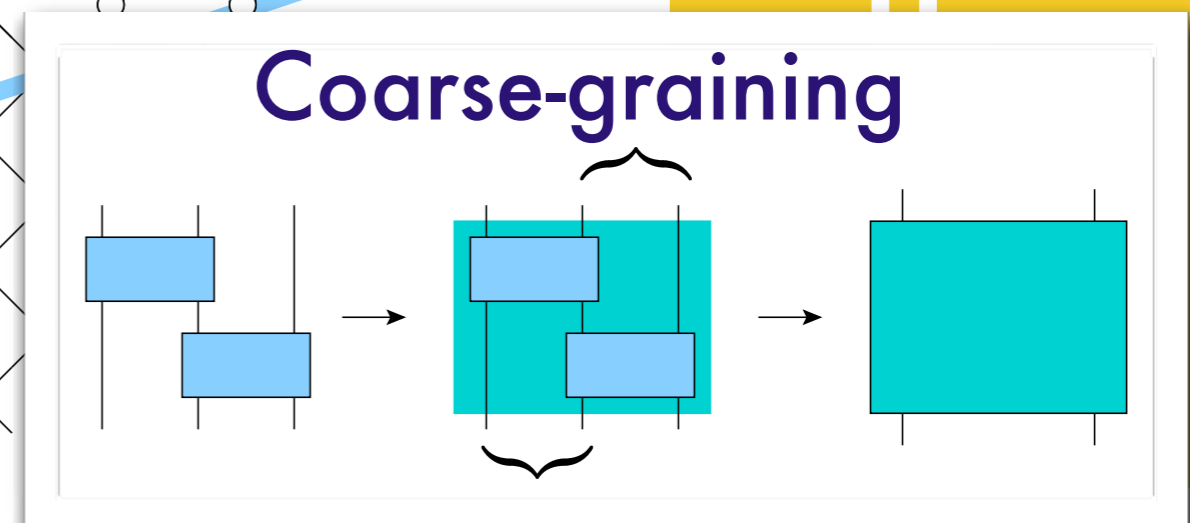
Simultaneity

Lorentz transformations from causality and topological homogeneity

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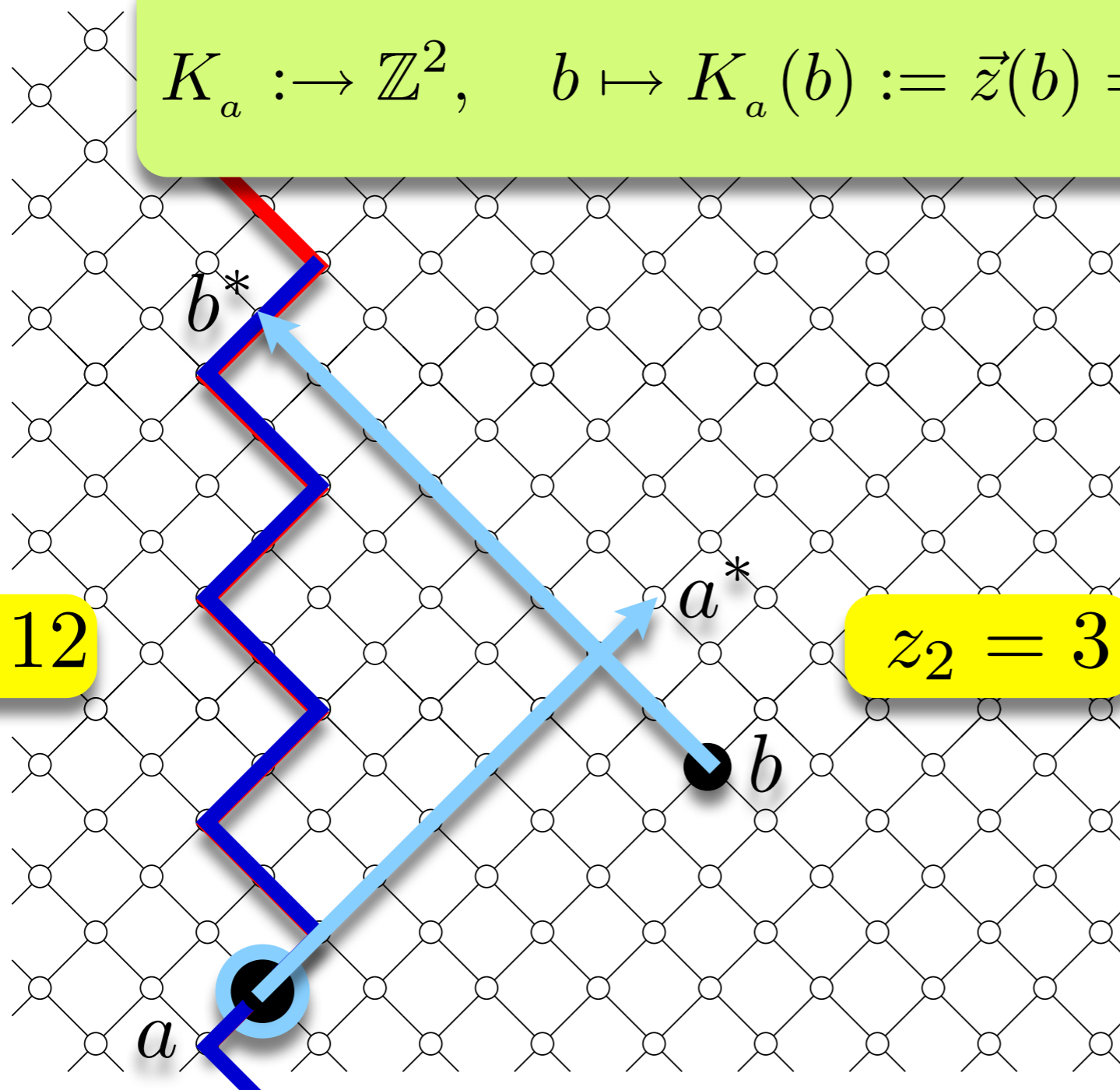
foliation



Lorentz transformations from causality and topological homogeneity

GMD and
A. Tosini
1008.4805

$$K_a : \rightarrow \mathbb{Z}^2, \quad b \mapsto K_a(b) := \vec{z}(b) = \begin{bmatrix} z_1(b) \\ z_2(b) \end{bmatrix}$$



Coordinates

Lorentz transformations from causality and topological homogeneity

GMD and
A. Tosini
1008.4805

Lemma 1 *An event $b \in L(O_a)$ belongs to the t -th leaf $L_t(O_a)$ for $t = (z_1 - z_2)/2$, and the number of events on such leaf between b and O_a is given by $s = (z_1 + z_2)/2$.*

According to the last Lemma the coordinates

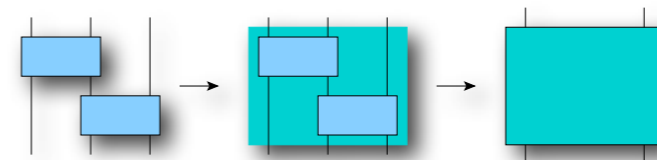
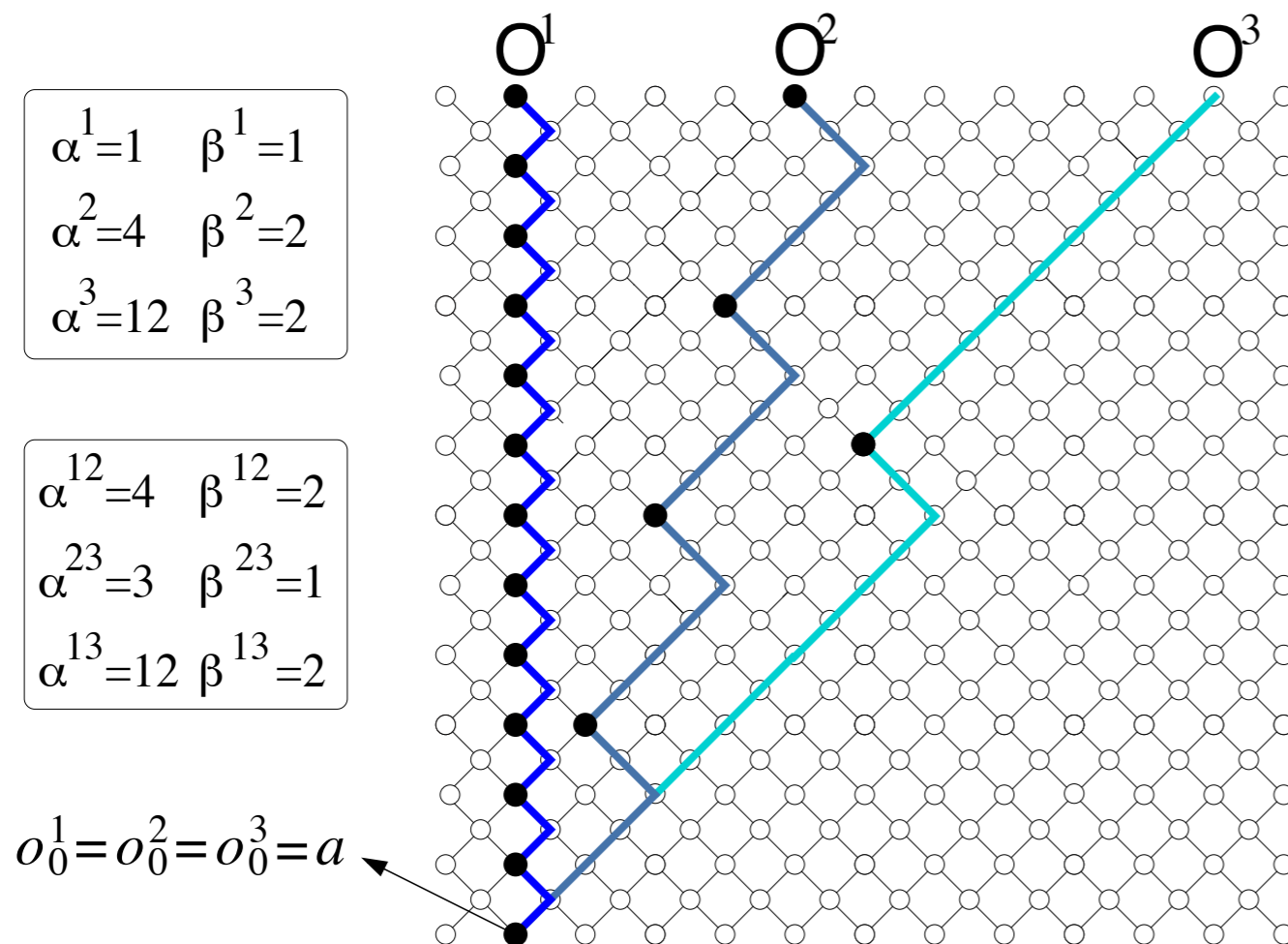
$$\begin{bmatrix} t(b) \\ s(b) \end{bmatrix} := 2^{\frac{1}{2}} \mathbf{U}(\pi/4) \begin{bmatrix} z(b) \\ z(b) \end{bmatrix}, \quad (10)$$

where $\mathbf{U}(\theta)$ is the matrix performing a θ -rotation, can be interpreted as the space-time coordinates of the event b in the frame $L(O_a)$.

Coordinates

Frames in standard configuration (boosted). Consider now two observers $O_a^1 = \{o_i^1\}$ and $O_a^2 = \{o_j^2\}$ sharing the same origin (homogeneity guarantees the existence of observers sharing the origin). We will shortly denote the two frames as \mathfrak{R}^1 and \mathfrak{R}^2 , and the corresponding coordinate maps as K^1 and K^2 . We will say that the two frames \mathfrak{R}^1 and \mathfrak{R}^2 are in *standard configuration* if there exist positive α^{12}, β^{12} , such that $\forall i \in \mathbb{Z}$

$$K^1(o_i^2) = \mathbf{D}^{12} K^2(o_i^2), \quad \mathbf{D}^{12} := \text{diag}(\alpha^{12}, \beta^{12}). \quad (11)$$

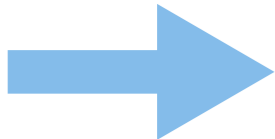


$$v^{12} = \frac{\alpha^{12} - \beta^{12}}{\alpha^{12} + \beta^{12}}$$

Boosts

Lorentz transformations from causality and topological homogeneity

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1008.4805


$$v^{13} = \frac{\alpha^{12}\alpha^{23} - \beta^{12}\beta^{23}}{\alpha^{12}\alpha^{23} + \beta^{12}\beta^{23}} = \frac{v_{12} + v_{23}}{1 + v_{12}v_{23}}$$



$$t^1 = \chi_{12} \frac{t^2 + v^{12}s^2}{\sqrt{1 - (v^{12})^2}}, \quad s^1 = \chi_{12} \frac{s^2 + v^{12}t^2}{\sqrt{1 - (v^{12})^2}},$$

$$\chi_{12} := \sqrt{\alpha^{12}\beta^{12}}$$

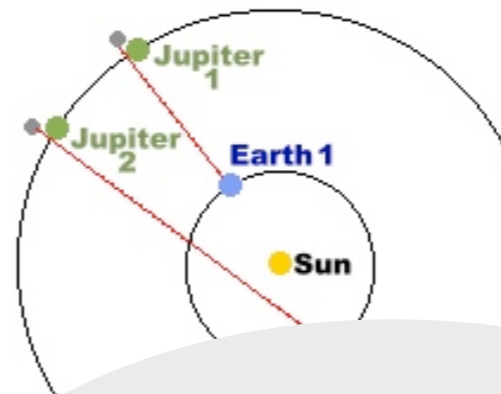
which differ from the Lorentz transformations only by the multiplicative factor χ_{12} . The factor χ_{12} can be removed by rescaling the coordinate map in Eq. (10) using the factor $(2\alpha\beta)^{\frac{1}{2}}$ in place of $2^{\frac{1}{2}}$, with the constants α and β

Coordinates

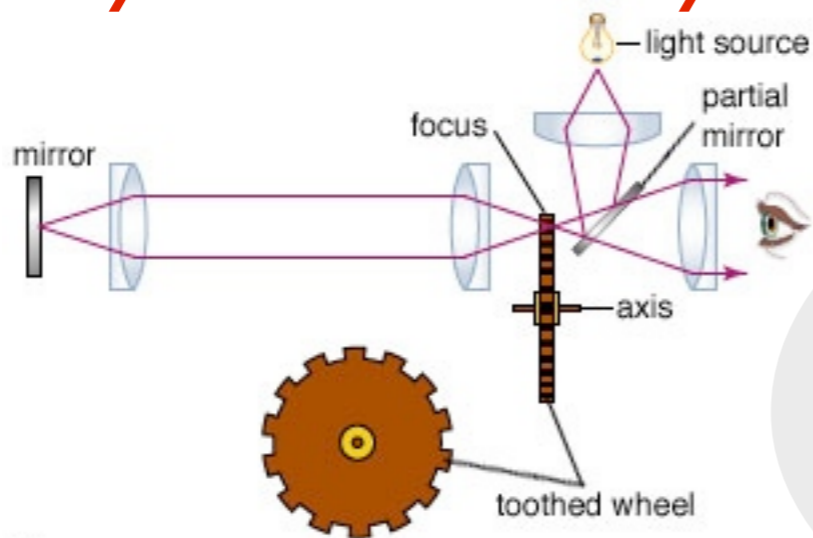
Conventionality of simultaneity, homogeneity, ...



The causal network manifests the *conventionality of simultaneity*.

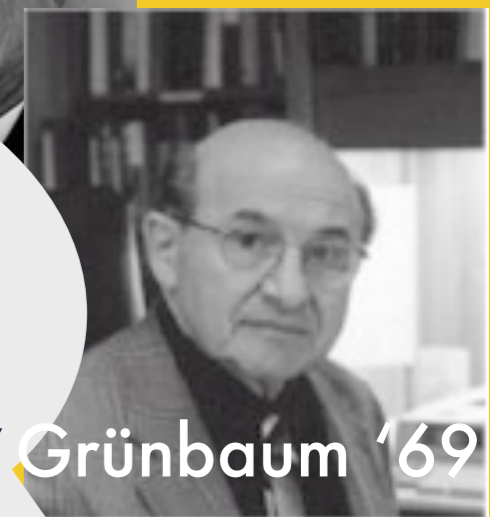


Bridgman '62



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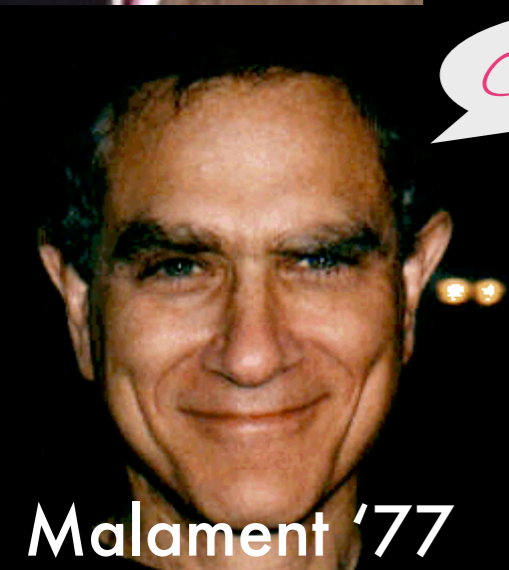
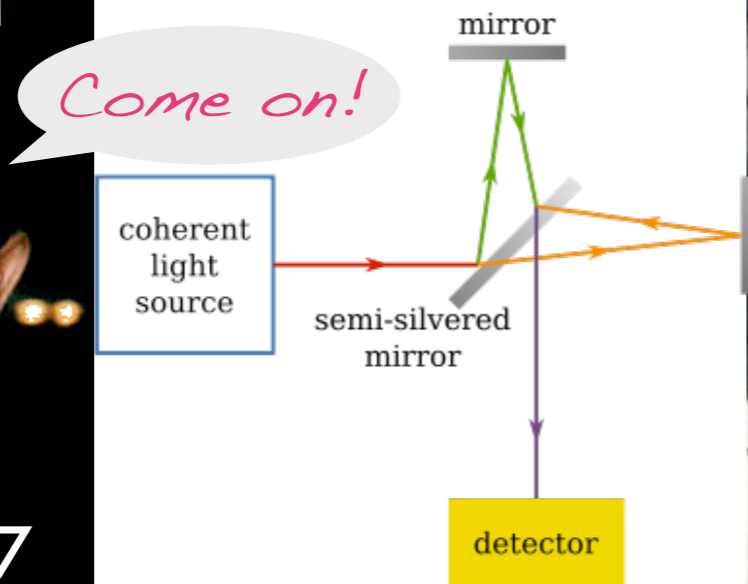
To determine simultaneity of distant events we need to know a speed, to measure a speed we need to know simultaneity of different events ... We can only determine the two-way average speed of light ...



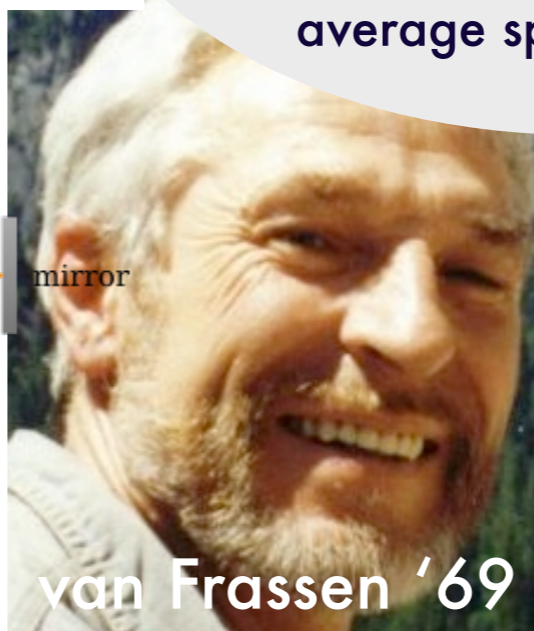
Grünbaum '69



Friedman '83



Malament '77



van Frassen '69



Salmon '69



Reichenback '57

WHAT IS THE INFORMATIONAL MEANING
OF INERTIAL MASS AND \hbar
AND HOW THE QUANTUM FIELD EMERGES

QC SIMULATION OF QFT

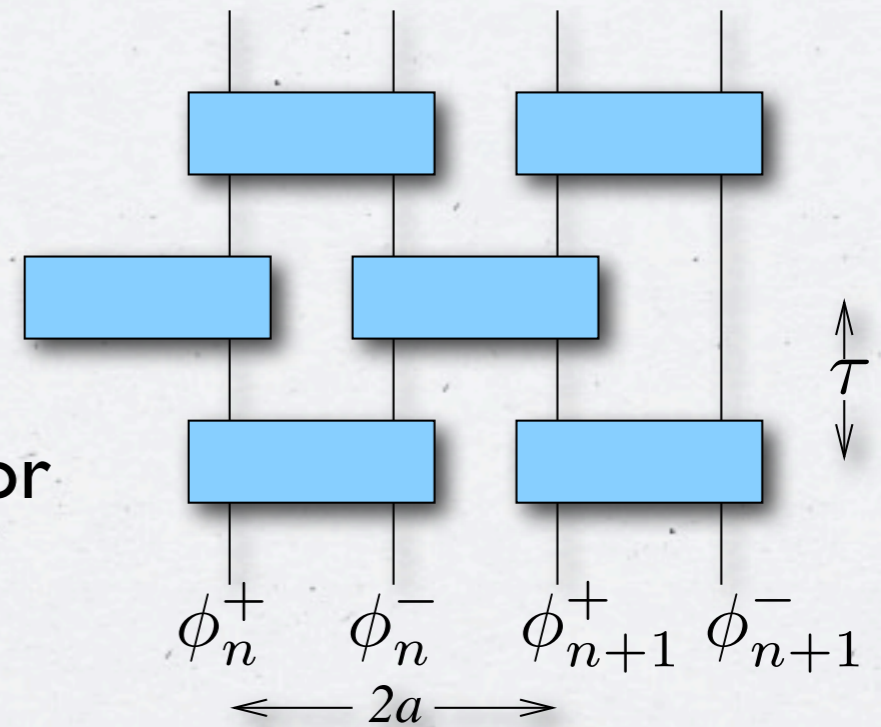
SIMPLE SCALAR FIELD IN 1 SPACE DIM.



a : **topon**: space-granularity (minimal in principle discrimination between independent events);

\mathcal{T} : **chronon**: time-granularity;

$\phi^\pm(x)$: (right/left propagating) field modes, operator function of space (evolving in time); we will describe it by the set of operators $\phi_n^\pm = a^{\frac{1}{2}} \phi^\pm(na)$



ϕ_n^\pm generally nonlocal operators. In QFT they satisfy (anti)commutation relations

Microcausality (equal time)

$$[\phi_n^\alpha, \phi_m^\beta]_\pm = \delta_{\alpha\beta} \delta_{nm}$$

+ : Fermi, - : Bose (Newton-Wigner)

CAUSAL SPEED


$$v_c := \frac{a}{\mathcal{T}}$$

QC SIMULATION OF QFT

VIOLATION OF EINSTEIN'S CAUSALITY

Simulation of QFT with a quantum computer, with gates performing infinitesimal transformations:

the simulation gives back exactly QFT in the limit $\tau, a \rightarrow 0$ and for infinite circuit, but ...



$v_c = \infty!$
Galileo!

Einstein causality only in average!

Lorentz-covariance cannot be derived from QT causality!

QC SIMULATION OF QFT

SIMPLE SCALAR FIELD IN 1 SPACE DIM.



Finite gate-transformations (not infinitesimal!)

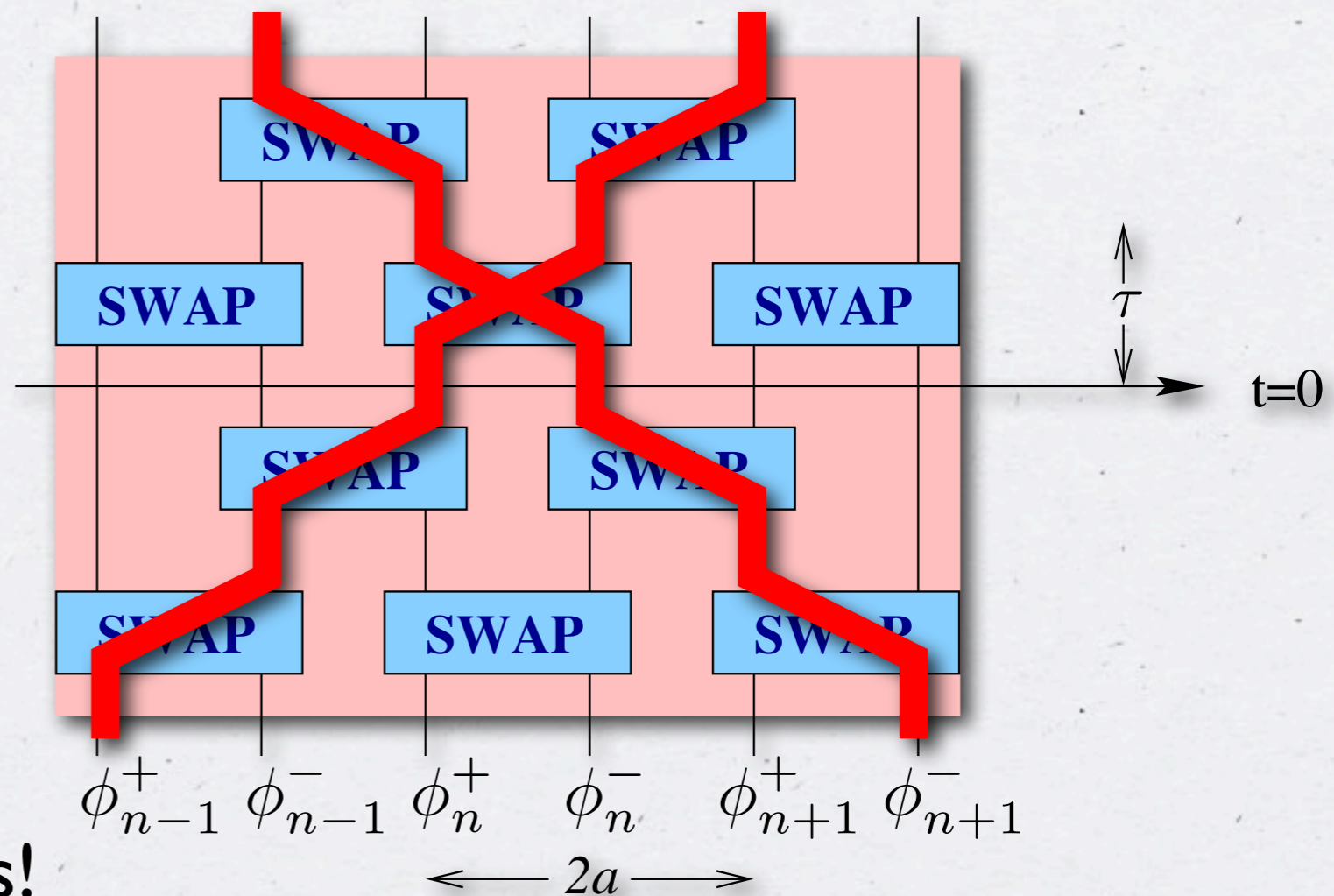
The causal speed v_c is finite!

Lorentz's transformations
emerge from the causal
network



Different QFT

observational consequences!



QC SIMULATION OF QFT

SIMPLE SCALAR FIELD IN 1 SPACE DIM.



Each gate evolves the field linearly:

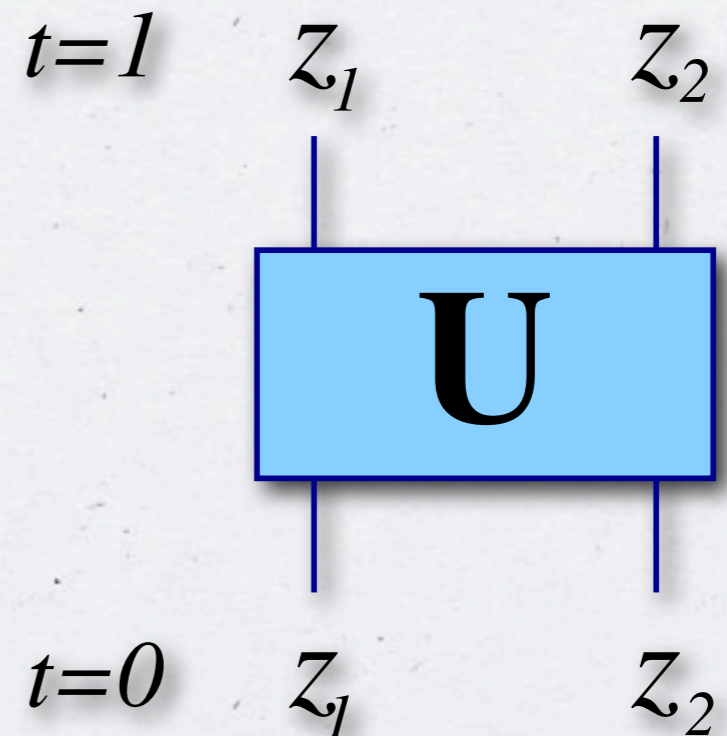
$$[z_i, z_j^\dagger]_{\pm} = \delta_{ij}$$

$$U z_n U^\dagger = \sum_k \mathbf{U}_{nk} z_k$$

$$\mathbf{U} = \|\mathbf{U}_{ij}\| \text{ unitary matrix}$$

Evolution from bipartite gates:

$$\begin{bmatrix} z_1 \\ z_2 \end{bmatrix}_{t=1} = U \begin{bmatrix} z_1 \\ z_2 \end{bmatrix} U^\dagger = \mathbf{U} \begin{bmatrix} z_1 \\ z_2 \end{bmatrix}$$



QC SIMULATION OF QFT

SIMPLE SCALAR FIELD IN 1 SPACE DIM.



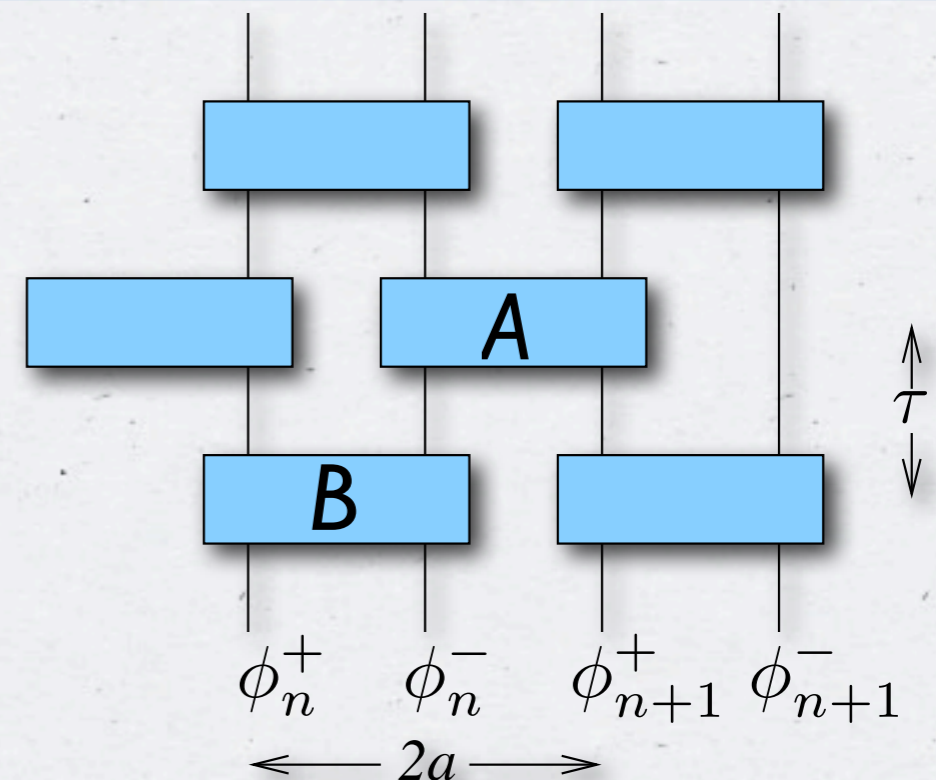
Commuting	Anticommuting
<i>Field is a local operator</i>	<i>Clifford algebraic construction</i>
<i>gates act on local algebras</i>	$\phi_n^+ = \sigma_{2n}^- \prod_{k=-\infty}^{n-1} \sigma_{2k+1}^z \sigma_{2k}^z$
	$\phi_n^- = \sigma_{2n+1}^- \sigma_{2n}^z \prod_{k=-\infty}^{n-1} \sigma_{2k+1}^z \sigma_{2k}^z$

According to Jordan-Schwinger:

$$A = A(\vec{\sigma}_{2n+1}, \vec{\sigma}_{2n+2})$$

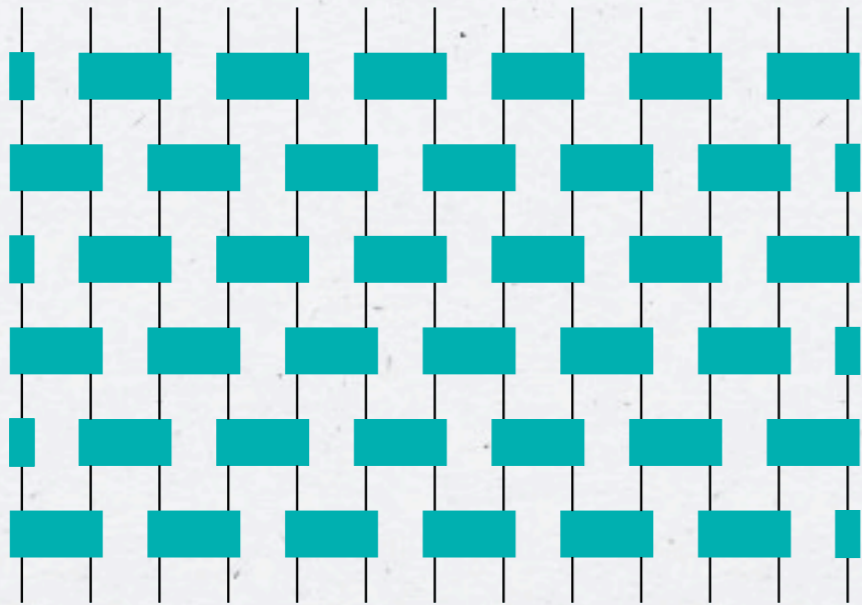
$$B = B(\vec{\sigma}_{2n}, \vec{\sigma}_{2n+1})$$

*i.e. gates act on local algebras
also for anticommuting fields*



THE NEW QCFT

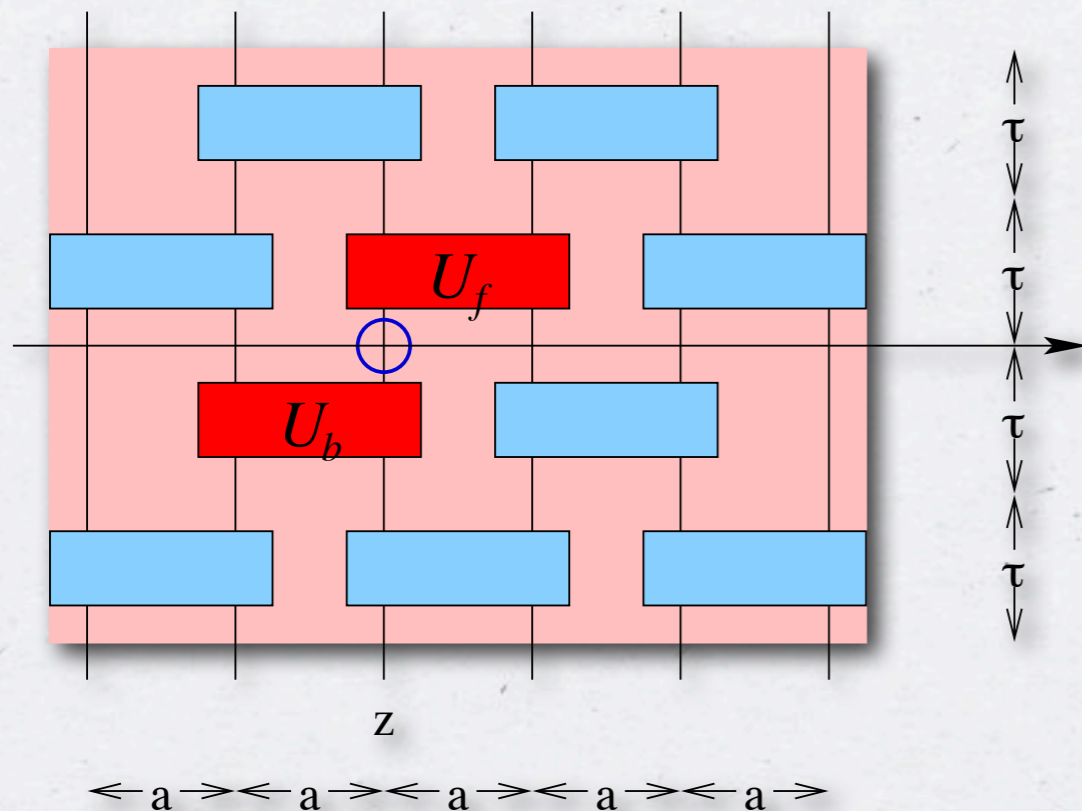
SIMPLE SCALAR FIELD IN 1 SPACE DIM.



Coarse-grained discrete derivatives:

$$\hat{\partial}_t z = \frac{1}{2k\tau} [z(k\tau) - z(-k\tau)]$$

$$\hat{\partial}_x = \frac{1}{4ka} (\delta_+^k - \delta_-^k)$$



“HAMILTONIAN”

$$H_{\text{gate}}^{(2n)} z = \frac{i}{2n\tau} [z(n\tau) - z(-n\tau)] = i\hat{\partial}_t z$$

$$H_{\text{gate}}^{(2)} z = \frac{i}{2\tau} (U_f z U_f^\dagger - U_b^\dagger z U_b)$$

THE NEW QCFT

MASSLESS FIELD

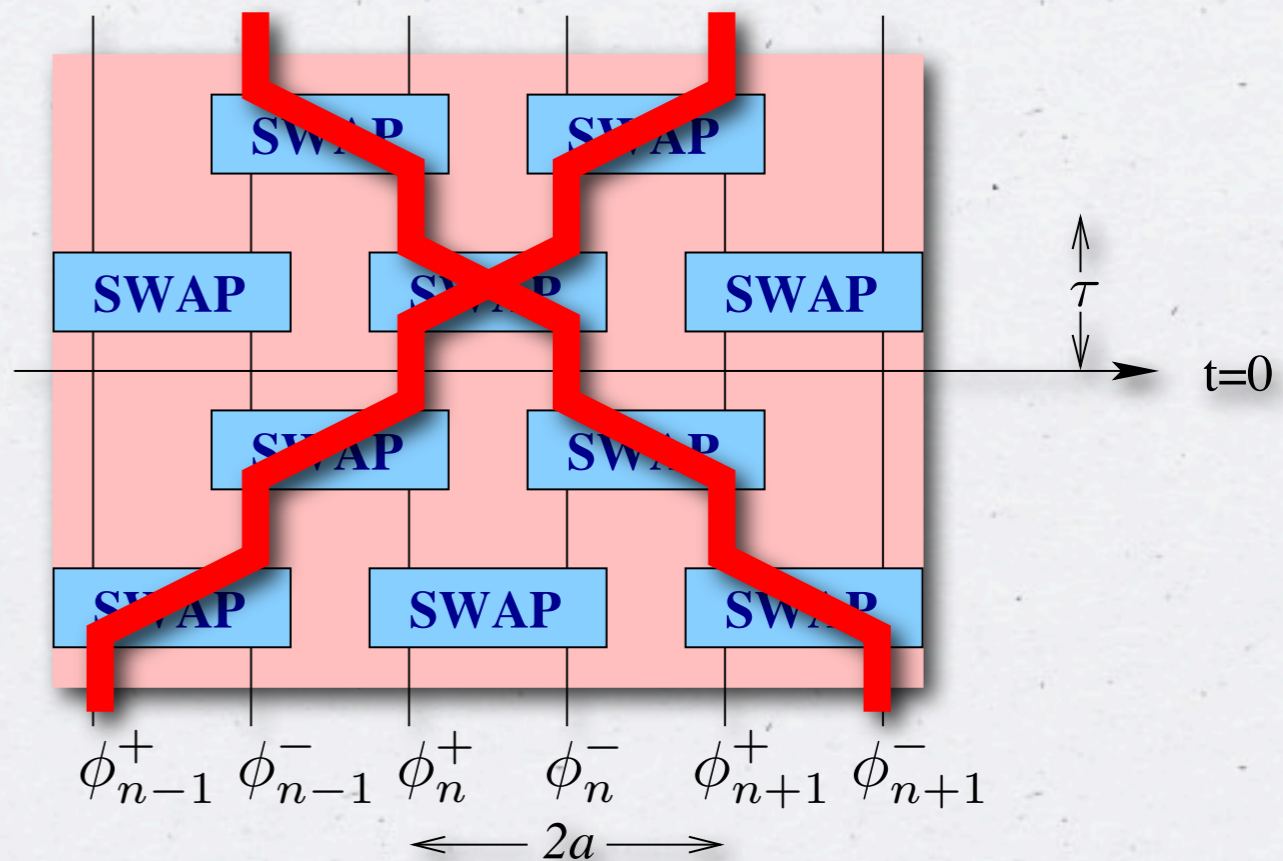


NEW QFT: *finite gate-transformations (not infinitesimal!)*

$$\phi_n^+(\pm 2\tau) = \phi_{n\pm 1}^+(0), \quad \phi_n^-(\pm 2\tau) = \phi_{n\mp 1}^-(0)$$

$$H_{\text{gate}}^{(4)} \phi_n^\alpha = i\alpha v_c \hat{\partial}_x \phi_n^\alpha, \quad \alpha = \pm$$

$$H_{\text{gate}}^{(4)} \phi_n = i\sigma_z v_c \hat{\partial}_x \phi_n$$



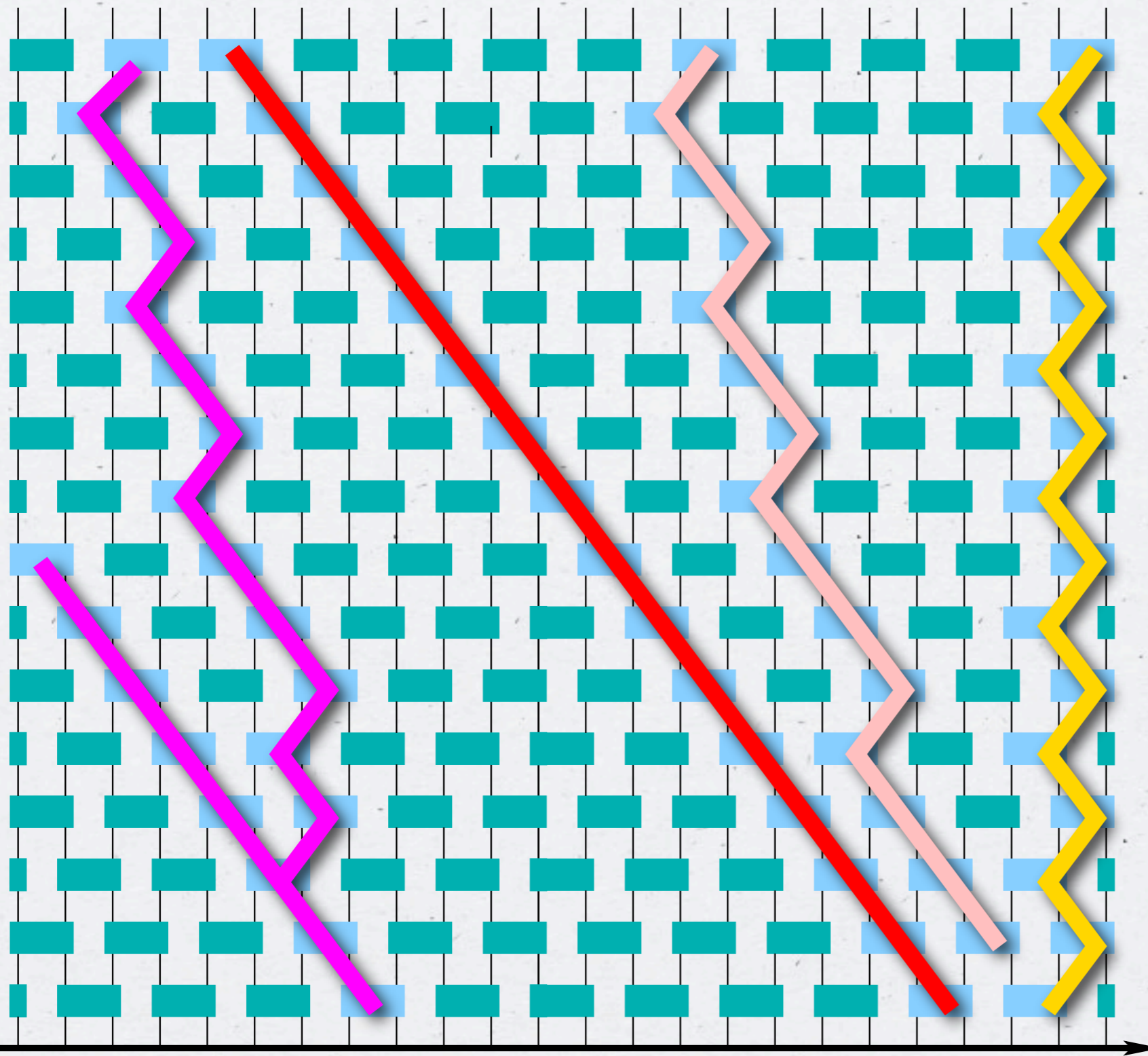
$$\hat{\square} \phi_n = 0$$

$$\hat{\square} = \hat{\partial}_x^2 - \frac{1}{v_c^2} \hat{\partial}_t^2$$

OBSERVATIONAL CONSEQUENCES: MASS-DEPENDENT REFRACTION INDEX OF VACUUM

THE NEW QCFT

WHAT IS INERTIAL MASS?



Zitterbewegungs

THE NEW QCFT

WHAT IS \hbar ?

$$m = \frac{\tau^2}{a^2} \hbar \nu$$

$$i\hat{\partial}_t \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix} = \begin{bmatrix} i\nu\hat{\partial}_x & \nu \\ \nu & -i\nu\hat{\partial}_x \end{bmatrix} \begin{bmatrix} \phi^+ \\ \phi^- \end{bmatrix}$$

m	mass in grams
a	topon
τ	chronon
ν	coupling between \pm modes
$\lambda = \frac{\hbar}{mc}$	Compton wavelenght

THE NEW QCFT

MASS-DEPENDENT VACUUM REFRACTION INDEX



Overall we must have: $U_f \phi_n^+ U_f - U_b^\dagger \phi_n^+ U_b = \zeta (\phi_{n+1}^+ - \phi_{n-1}^+) - 4i \frac{a}{\lambda} \phi_n^-$

For local gates involving only n.n. wires, the overall forward and backward unitary interactions involving a minimal number of field operators have the form

$$U_f \phi_n^+ U_f^\dagger = \eta \phi_n^+ + \zeta \phi_{n+1}^+ + \gamma \phi_n^- \quad U_b^\dagger \phi_n^+ U_b = \eta \phi_n^+ + \zeta \phi_{n-1}^+ + \gamma' \phi_n^-$$

with $\zeta > 0$ and $\gamma - \gamma' = -4i \frac{a}{\lambda}$

But normalization of the row of the unitary matrix corresponds to:

$$|\gamma|, |\gamma'| \leq \sqrt{1 - \zeta^2} \implies \frac{2a}{\lambda} \leq \sqrt{1 - \zeta^2}$$

which is the bound for the vacuum refraction index:

$$\zeta^{-1} \geq \left[1 - \left(\frac{2a}{\lambda} \right)^2 \right]^{-\frac{1}{2}}$$

Effect visible at large mass $\lambda \sim 2a$

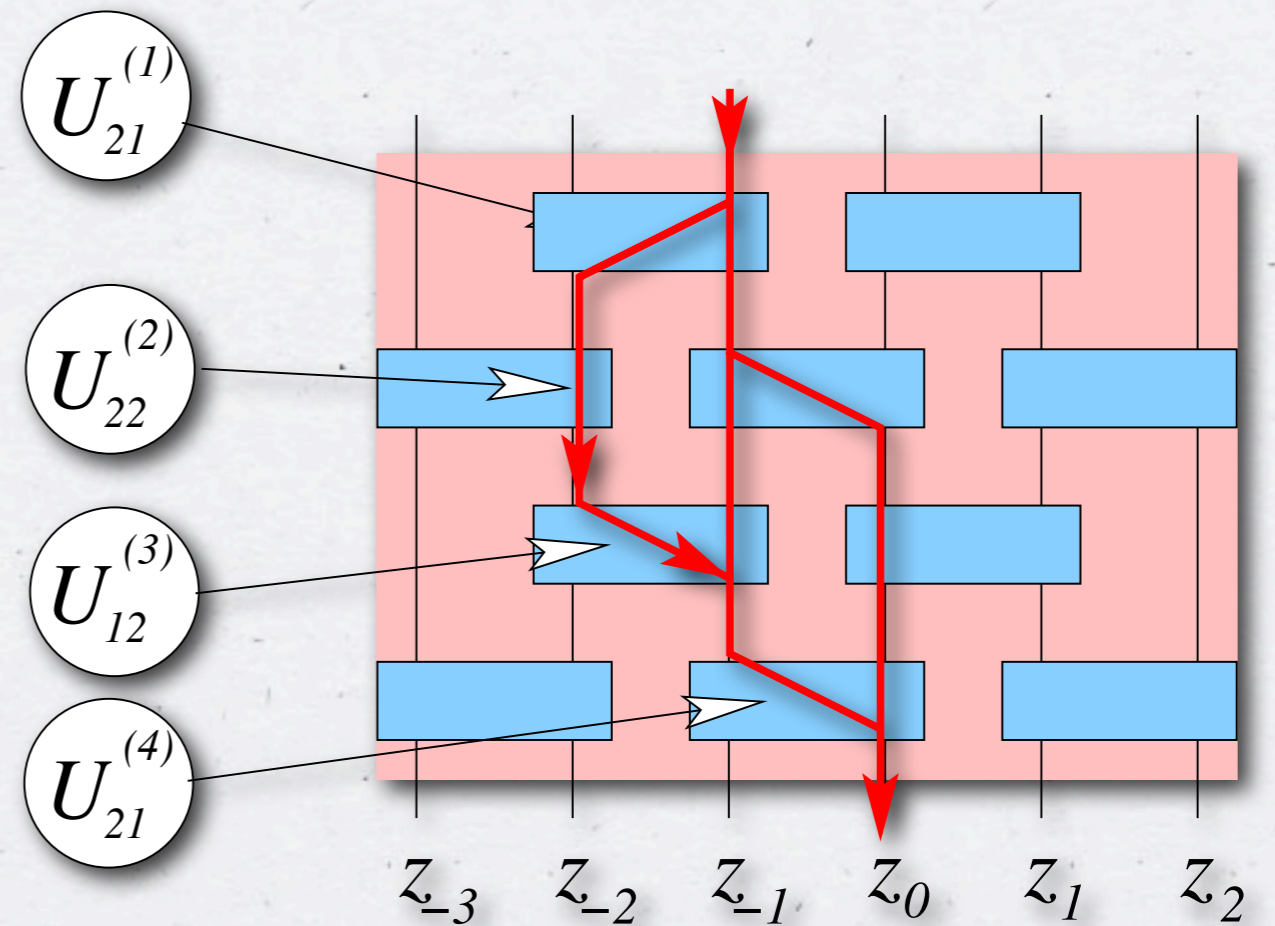
THE NEW QCFT

FEYNMAN PATH-SUM



We need to develop a *path-sum calculus* over the circuit:

1. Number all the input wires at each gate, from the leftmost to the rightmost one, and do the same for the output wires
2. We say that a wire l is in the past-cone of the wire k if there is a path from l to k passing through gates.
3. For any output wire k and any input wire l in its causal past cone, consider all paths connecting k with l
4. The following linear expansion holds



$$z_l(t) = \sum_{\mathbf{i}_{kl}} U_{i_1 i_2}^{(1)} U_{i_2 i_3}^{(2)} \cdots U_{i_n i_{n+1}}^{(n)} z_k(0)$$

$$\mathbf{i}_{kl} = (i_1 i_2 \dots i_n i_{n+1}) \text{ with } i_1 = k, i_{n+1} = l,$$

THE NEW QCFT

KLEIN GORDON WITH MASS

“Hamiltonian”

$$H_{\text{gate}}^{(4)} = \frac{i}{4\tau} \begin{bmatrix} A_{21}B_{21}\delta_- - B_{12}^\dagger A_{12}^\dagger \delta_+ + A_{22}B_{11} - B_{11}^\dagger A_{22}^\dagger & (A_{21}B_{22} - B_{11}^\dagger A_{21}^\dagger)\delta_- + A_{22}B_{12} - B_{12}^\dagger A_{11}^\dagger \\ (A_{12}B_{11} - B_{22}^\dagger A_{12}^\dagger)\delta_+ + A_{11}B_{21} - B_{21}^\dagger A_{22}^\dagger & A_{12}B_{12}\delta_+ - B_{21}^\dagger A_{21}^\dagger \delta_- + A_{11}B_{22} - B_{22}^\dagger A_{11}^\dagger \end{bmatrix}$$

Hermiticity is satisfied:

$$\langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle = \langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle^* \implies i(A_{aa}B_{bb} - A_{aa}^\dagger B_{bb}^\dagger) \in \mathbb{R},$$

$$\langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_n^\mp \rangle = \langle \phi_n^\mp | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle^* \implies (A_{22}B_{12} - A_{11}^\dagger B_{12}^\dagger) = -(A_{11}B_{21} - A_{22}^\dagger B_{21}^\dagger)^*,$$

$$\langle \phi_{n+1}^\pm | H_{\text{gate}}^{(4)} | \phi_n^\pm \rangle = \langle \phi_n^\pm | H_{\text{gate}}^{(4)} | \phi_{n+1}^\pm \rangle^* \implies A_{ab}^\dagger B_{ab}^\dagger = A_{ba}^* B_{ba}^*,$$

$$\langle \phi_n^+ | H_{\text{gate}}^{(4)} | \phi_{n-1}^- \rangle = \langle \phi_n^- | H_{\text{gate}}^{(4)} | \phi_{n+1}^+ \rangle^* \implies A_{21}B_{22} - A_{21}^\dagger B_{11}^\dagger = -(A_{12}B_{11} - A_{12}^\dagger B_{22}^\dagger)^*.$$

THE NEW QCFT

KLEIN GORDON WITH MASS

Using the smoothness constraint $\frac{1}{2}(\delta_+ + \delta_-) \simeq 1$ corresponding to $\delta_{\pm} = 1 \pm 2a\hat{\partial}_x$ one writes the Hamiltonian in the Dirac fashion:

$$H_{\text{gate}}^{(4)} = v_c(\mathbf{H} + \mathbf{K}\hat{\partial}_x), \quad \mathbf{H} := \begin{bmatrix} H_{11} & H_{12} \\ H_{12}^* & H_{22} \end{bmatrix}, \quad \mathbf{K} := \begin{bmatrix} K_{11} & K_{12} \\ -K_{12}^* & K_{22} \end{bmatrix},$$

Restrict to KG with propagation speed ζc , namely

$$(H_{\text{gate}}^{(4)})^2 = -c^2(\zeta^2 \hat{\partial}_x^2 - \lambda^{-2})$$

$$\lambda := \frac{\hbar}{mc} = 3.86159 * 10^{-13} \text{m}$$

Compton wavelength

After some manipulations one gets the **mass-dependent vacuum refraction index**:

$$\zeta^{-1} = \frac{2}{1 + \sqrt{1 - \left(\frac{2a}{\lambda}\right)^2}}$$

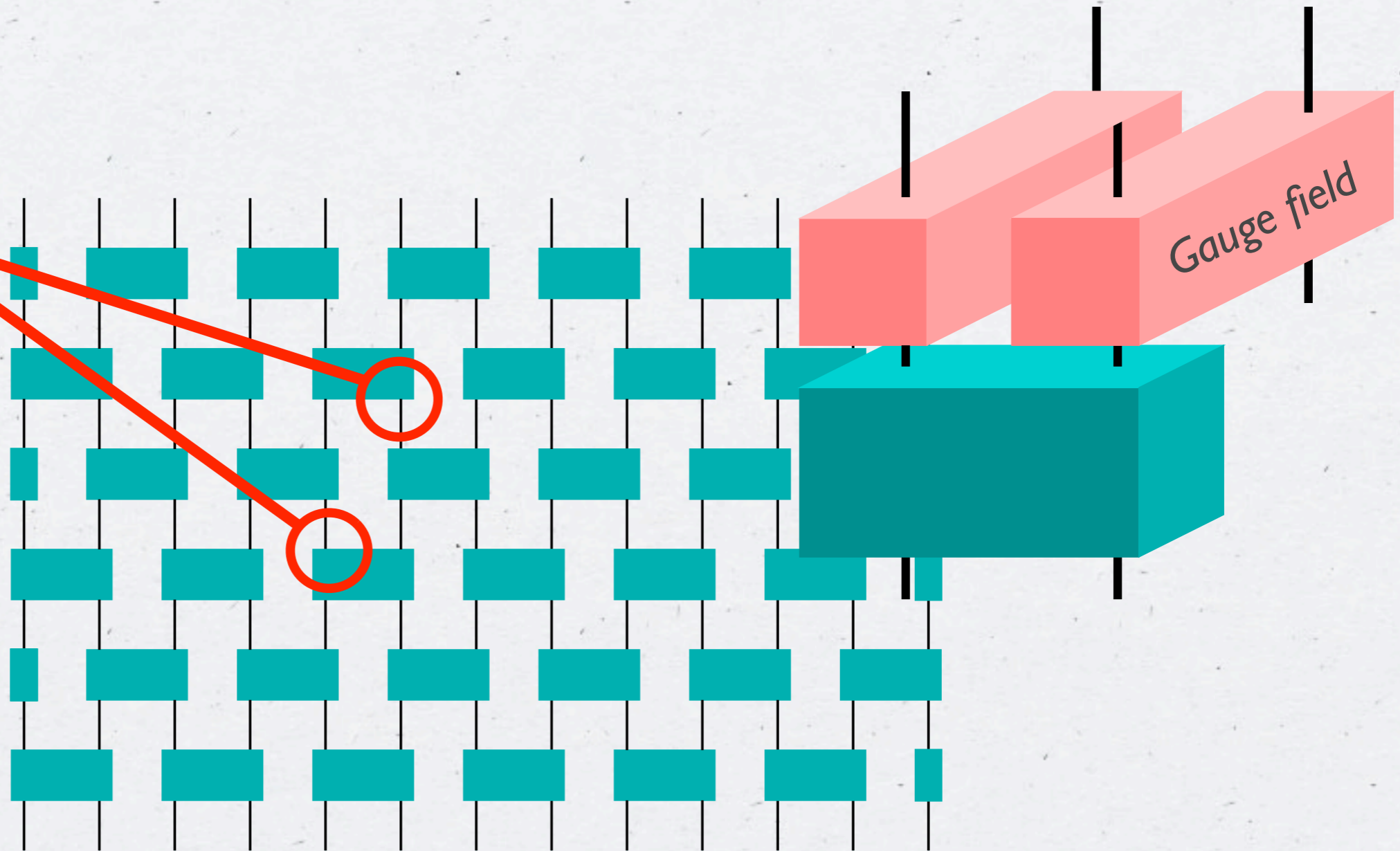
Easy to generalize to particles with spin

THE NEW QCFT

GAUGE INVARIANCE



$$U(x)$$



***Natively nonabelian Gauge theory!
and on ... foliation !!!***



***Good for
QGravity?***

THE NEW QCFT

QCFT FOR MORE THAN 1 SPACE DIM?



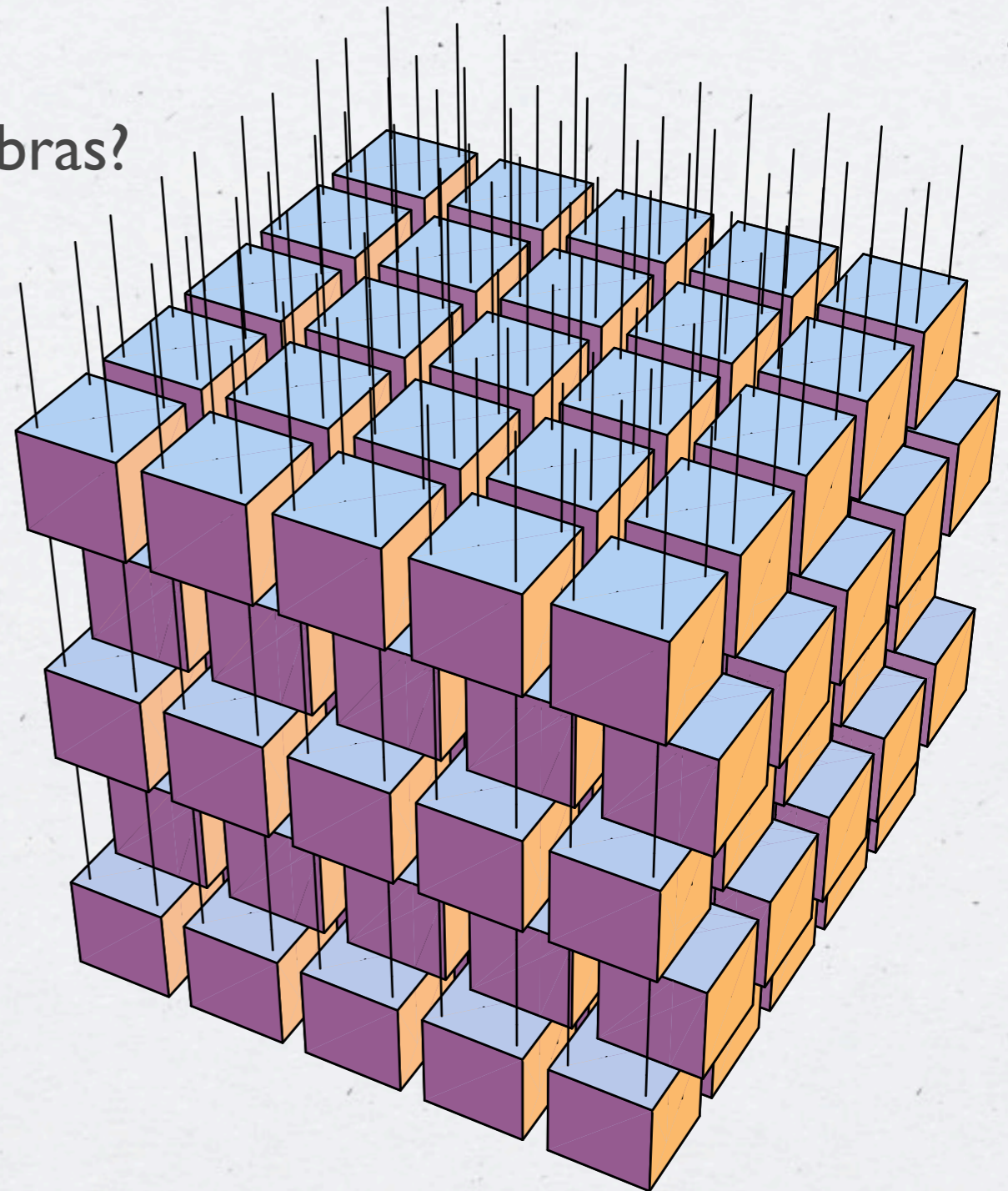
Need six space field operators

Anticommuting fields in terms of local algebras?

Do we really need anticommuting fields?

Grassman variables?

Microcausality and parastatistics



THE NEW QCFT

COMPARED TO THE USUAL QFT



QCFT solve many problems that plague QFT:

- * Feynman's path integral
- * u.v. renormalization
- * no need of quantization rules (must be emergent)
- * problems related to the continuous
- * ...

The idea is to regard the QCFT as the “true theory”, and the usual QFT as an approximation for “mesoscopic” scale

CONCLUSIONS

PHYSICS IS INFORMATION

- * Quantum Theory is an information theory
- * Space-time and relativistic covariance emerge from the information processing
- * The whole Physics is emergent (inertial mass, Planck constant, quantization rules, ...)
- * The new causal QCFT:
 - * has no space-background (QG-ready)
 - * doesn't need quantization
 - * cures many problems that plague QFT
 - * opens a new route to foundations of QFT
 - * has empirical consequences ...

TODO SOON

PHYSICS IS INFORMATION



- * Quantization rules as emergent
- * General correspondence: Lagrangian-gates
- * Anticommuting fields and (parastatistics)
- * Emergent unitary representation of the Lorentz group
- * Violations of Lorentz-covariance
- * Informational meaning of energy, gravitational mass,...
- * Build up a complete QCFT for Dirac in 3 space dimensions
- * Experimental consequences ...

THANK YOU
